Periodic Homogenization of Green and Neumann Functions

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Abstract

For a family of second-order elliptic operators with rapidly oscillating periodic coefficients, we study the asymptotic behavior of the Green and Neumann functions, using Dirichlet and Neumann correctors. As a result we obtain asymptotic expansions of Poisson kernels and the Dirichlet-to-Neumann maps as well as near optimal convergence rates in $W^{1,p}$ for solutions with Dirichlet or Neumann boundary conditions.

1 Introduction

The main purpose of this paper is to study the asymptotic behavior of the Green and Neumann functions for a family of elliptic operators with rapidly oscillating coefficients. More precisely, consider

$$\mathcal{L}_{\varepsilon} = -\operatorname{div}(A(x/\varepsilon)\nabla) = -\frac{\partial}{\partial x_i} \left[a_{ij}^{\alpha\beta} \left(\frac{x}{\varepsilon} \right) \frac{\partial}{\partial x_j} \right], \quad \varepsilon > 0$$
 (1.1)

(the summation convention is used throughout the paper). We will assume that $A(y) = (a_{ij}^{\alpha\beta}(y))$ with $1 \le i, j \le d$ and $1 \le \alpha, \beta \le m$ is real and satisfies the ellipticity condition

$$\mu|\xi|^2 \le a_{ij}^{\alpha\beta}(y)\xi_i^{\alpha}\xi_j^{\beta} \le \frac{1}{\mu}|\xi|^2 \quad \text{for } y \in \mathbb{R}^d \text{ and } \xi = (\xi_i^{\alpha}) \in \mathbb{R}^{dm}, \tag{1.2}$$

where $\mu > 0$, and the periodicity condition

$$A(y+z) = A(y)$$
 for $y \in \mathbb{R}^d$ and $z \in \mathbb{Z}^d$. (1.3)

We will also impose the smoothness condition

$$|A(x) - A(y)| \le \tau |x - y|^{\lambda}$$
 for some $\lambda \in (0, 1]$ and $\tau \ge 0$. (1.4)

Let $G_{\varepsilon}(x,y) = (G_{\varepsilon}^{\alpha\beta}(x,y))$ and $N_{\varepsilon}(x,y) = (N_{\varepsilon}^{\alpha\beta}(x,y))$ denote the Green and Neumann functions respectively, for $\mathcal{L}_{\varepsilon}$ in a bounded domain Ω , with pole at y. We are interested in

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the asymptotic behavior, as $\varepsilon \to 0$, of $G_{\varepsilon}(x,y)$, $N_{\varepsilon}(x,y)$, $\nabla_x G_{\varepsilon}(x,y)$ and $\nabla_x N_{\varepsilon}(x,y)$ as well as $\nabla_x \nabla_y G_{\varepsilon}(x,y)$ an $\nabla_x \nabla_y N_{\varepsilon}(x,y)$. We shall use $G_0(x,y)$ and $N_0(x,y)$ to denote the Green and Neumann functions respectively, for the homogenized (effective) operator \mathcal{L}_0 in Ω .

Let $P_j^{\beta} = x_j(0,\ldots,1,\ldots)$ with 1 in the β^{th} position for $1 \leq j \leq d$ and $1 \leq \beta \leq m$. To state our main results, we need to introduce the matrix of Dirichlet correctors $\Phi_{\varepsilon,j}^{\beta} = (\Phi_{\varepsilon,j}^{1\beta},\ldots,\Phi_{\varepsilon,j}^{m\beta})$ in Ω , defined by

$$\mathcal{L}_{\varepsilon}(\Phi_{\varepsilon,j}^{\beta}) = 0 \quad \text{in } \Omega \quad \text{and} \quad \Phi_{\varepsilon,j}^{\beta} = P_{j}^{\beta} \quad \text{on } \partial\Omega,$$
 (1.5)

as well as the matrix of Neumann correctors $\Psi_{\varepsilon,j}^{\beta} = \left(\Psi_{\varepsilon,j}^{1\beta}, \dots, \Psi_{\varepsilon,j}^{m\beta}\right)$ in Ω , defined by

$$\mathcal{L}_{\varepsilon}(\Psi_{\varepsilon,j}^{\beta}) = 0 \quad \text{in } \Omega \quad \text{and} \quad \frac{\partial}{\partial \nu_{\varepsilon}} (\Psi_{\varepsilon,j}^{\beta}) = \frac{\partial}{\partial \nu_{0}} (P_{j}^{\beta}) \quad \text{on } \partial \Omega.$$
 (1.6)

Here $\partial/\partial\nu_{\varepsilon}$ denotes the conormal derivative associated with $\mathcal{L}_{\varepsilon}$ for $\varepsilon \geq 0$.

The following are the main results of the paper.

Theorem 1.1. Let $\mathcal{L}_{\varepsilon} = -\text{div}(A(x/\varepsilon)\nabla)$ with the matrix A(y) satisfying conditions (1.2), (1.3) and (1.4). Then for any $x, y \in \Omega$,

$$|G_{\varepsilon}(x,y) - G_0(x,y)| \le \frac{C\varepsilon}{|x-y|^{d-1}}$$
(1.7)

if Ω is a bounded $C^{1,1}$ domain, and

$$\left| \frac{\partial}{\partial x_i} \left\{ G_{\varepsilon}^{\alpha\beta}(x, y) \right\} - \frac{\partial}{\partial x_i} \left\{ \Phi_{\varepsilon, j}^{\alpha\gamma}(x) \right\} \cdot \frac{\partial}{\partial x_j} \left\{ G_0^{\gamma\beta}(x, y) \right\} \right| \le \frac{C\varepsilon \ln\left[\varepsilon^{-1}|x - y| + 2\right]}{|x - y|^d} \tag{1.8}$$

if Ω is a bounded $C^{2,\eta}$ domain for some $\eta \in (0,1)$, where C depends only on d, m, μ , τ , λ and Ω .

Theorem 1.2. Suppose that A(y) satisfies the same conditions as in Theorem 1.1. Also assume that $A^* = A$, i.e., $a_{ij}^{\alpha\beta}(y) = a_{ji}^{\beta\alpha}(y)$ for $1 \le i, j \le d$ and $1 \le \alpha, \beta \le m$. Then for any $x, y \in \Omega$,

$$|N_{\varepsilon}(x,y) - N_0(x,y)| \le \frac{C\varepsilon \ln\left[\varepsilon^{-1}|x-y|+2\right]}{|x-y|^{d-1}}$$
(1.9)

if Ω is a bounded $C^{1,1}$ domain, where C depends only on d, m, μ , τ , λ and Ω . Moreover, if Ω is a bounded $C^{2,\eta}$ domain for some $\eta \in (0,1)$,

$$\left| \frac{\partial}{\partial x_i} \left\{ N_{\varepsilon}^{\alpha\beta}(x, y) \right\} - \frac{\partial}{\partial x_i} \left\{ \Psi_{\varepsilon, j}^{\alpha\gamma}(x) \right\} \cdot \frac{\partial}{\partial x_j} \left\{ N_0^{\gamma\beta}(x, y) \right\} \right| \le \frac{C_t \, \varepsilon^t \ln \left[\varepsilon^{-1} M + 2 \right]}{|x - y|^{d - 1 + t}} \tag{1.10}$$

for any $x, y \in \Omega$ and $t \in (0,1)$, where $M = \operatorname{diam}(\Omega)$ and C_t depends only on d, m, μ , τ , λ , t and Ω .

A few remarks are in order.

Remark 1.3. In the case of a scalar equation (m = 1), the estimate (1.7) holds for bounded measurable coefficients satisfying (1.2) and (1.3) (see Theorem 3.3).

Remark 1.4. The matrix of Dirichlet correctors $(\Phi_{\varepsilon,j}^{\beta})$ was introduced in [2] to establish the boundary Lipschitz estimates for solutions with Dirichlet conditions, while the matrix of Neumann correctors $(\Psi_{\varepsilon,\beta}^{\beta})$ was introduced in [21] to establish the same estimates for solutions with Neumann boundary conditions. It is known that $\|\Phi_{\varepsilon,j}^{\beta} - P_j^{\beta}\|_{L^{\infty}(\Omega)} \leq C\varepsilon$ and $\|\nabla \Phi_{\varepsilon,j}^{\beta}\|_{L^{\infty}(\Omega)} + \|\nabla \Psi_{\varepsilon,j}^{\beta}\|_{L^{\infty}(\Omega)} \leq C\varepsilon$ Under the condition $\Psi_{\varepsilon,j}^{\beta}(x_0) = P_j^{\beta}(x_0)$ for some $x_0 \in \Omega$, we also have $\|\Psi_{\varepsilon,j}^{\beta} - P_j^{\beta}\|_{L^{\infty}(\Omega)} \leq C\varepsilon \ln[\varepsilon^{-1}M + 2]$ (see Propositions 2.4 and 2.5).

Remark 1.5. Estimates (1.7) and (1.9) in Theorems 1.1 and 1.2 allow us to establish $O(\varepsilon)$ estimates for $||u_{\varepsilon} - u_0||_{L^p(\Omega)}$ ($1) for solutions with Dirichlet or Neumann boundary conditions (see Theorems 3.4 and 4.5). More importantly, estimates (1.8) and (1.10) yield near optimal convergence rates in <math>W^{1,p}(\Omega)$ for any $1 . In fact, let <math>\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω and $u_{\varepsilon} = 0$ on $\partial\Omega$. Then

$$\|u_{\varepsilon} - u_{0} - \left\{\Phi_{\varepsilon,j}^{\beta} - P_{j}^{\beta}\right\} \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \|_{W_{0}^{1,p}(\Omega)} \le C_{p} \varepsilon \left\{\ln[\varepsilon^{-1}M + 2]\right\}^{4|\frac{1}{2} - \frac{1}{p}|} \|F\|_{L^{p}(\Omega)}. \tag{1.11}$$

In the case of Neumann boundary conditions we obtain

$$\|u_{\varepsilon} - u_0 - \{\Psi_{\varepsilon,j}^{\beta} - P_j^{\beta}\} \frac{\partial u_0^{\beta}}{\partial x_i} \|_{W^{1,p}(\Omega)} \le C_{t,p} \varepsilon^t \|F\|_{L^p(\Omega)}$$
(1.12)

for any $t \in (0,1)$, where $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω , $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = 0$ on $\partial \Omega$ and $\int_{\Omega} F = \int_{\partial \Omega} u_{\varepsilon} = 0$. See subsections 3.2 and 4.2 for details. Let $w_{\varepsilon} = u_{\varepsilon} - u_0 - \varepsilon \chi_j^{\beta}(x/\varepsilon) \frac{\partial u_0^{\beta}}{\partial x_j}$, where $(\chi_j^{\beta}(y))$ denotes the matrix of correctors for \mathcal{L}_1 in \mathbb{R}^d . Estimates (1.11) and (1.12) should be compared to the well known $O(\varepsilon^{1/2})$ estimate: $||w_{\varepsilon}||_{H^1(\Omega)} = O(\varepsilon^{1/2})$ (see e.g. [6]), and to the following estimate,

$$||w_{\varepsilon}||_{H^{1/2}(\Omega)} + \left\{ \int_{\Omega} |\nabla w_{\varepsilon}(x)|^2 \operatorname{dist}(x, \partial \Omega) \, dx \right\}^{1/2} \le C \varepsilon ||F||_{L^2(\Omega)}, \tag{1.13}$$

proved in [20, Theorems 3.4 and 5.2]. Due to the presence of a boundary layer, both the $O(\varepsilon^{1/2})$ estimate and (1.13) are more or less sharp. The Dirichlet and Neumann correctors are introduced precisely to deal with boundary layer phenomena in periodic homogenization.

Remark 1.6. Our approach to Theorems 1.1 and 1.2 also leads to asymptotic estimates of $\nabla_x \nabla_y G_{\varepsilon}(x,y)$ and $\nabla_x \nabla_y N_{\varepsilon}(x,y)$ (see subsection 3.3 and Remark 4.9). As a result we obtain asymptotic expansions for $(\partial/\partial x_i)(\mathcal{L}_{\varepsilon})^{-1}(\partial/\partial x_j)$ and Λ_{ε} , the Dirichlet-to-Neumann map associated with $\mathcal{L}_{\varepsilon}$.

The asymptotic expansion of the fundamental solutions as well as the heat kernels for $\mathcal{L}_{\varepsilon}$ in \mathbb{R}^d has been studied, using the method of Bloch waves; see e.g. [26, 11] and their references (also see [24] for results obtained by the method of G-convergence). In the presence of boundary, the Bloch representation is no longer available. In a series of papers [2, 4, 3], M. Avellaneda and F. Lin introduced the compactness methods, which originated in

the regularity theory in the calculus of variations and minimal surfaces, to homogenization problems. In particular, they established an asymptotic expansion for $\nabla_y G_{\varepsilon}(x,y)$ in [4], using Dirichlet correctors. As a result, it was proved in [4] that if Ω is $C^{1,\eta}$ for some $\eta \in (0,1)$,

$$P_{\varepsilon}(x,y) = P_0(x,y)\omega_{\varepsilon}(y) + R_{\varepsilon}(x,y)$$
 for $x \in \Omega$ and $y \in \partial\Omega$, (1.14)

where $P_{\varepsilon}(x,y)$ ($\varepsilon \geq 0$) denotes the Poisson kernel for $\mathcal{L}_{\varepsilon}$ in Ω ,

$$\omega_{\varepsilon}(y) = \frac{\partial}{\partial n(y)} \left\{ \Phi_{\varepsilon,k}^*(y) \right\} \cdot n_k(y) h(y) \cdot n_i(y) n_j(y) a_{ij}(y/\varepsilon) \tag{1.15}$$

and the remainder $R_{\varepsilon}(x,y)$ satisfies

$$\lim_{\varepsilon \to 0} \sup \left\{ |R_{\varepsilon}(x, y)| : \ x \in E \text{ and } y \in \partial \Omega \right\} = 0$$
 (1.16)

for any compact subset E of Ω (the results are stated for the case m=1; however, the argument in [4] works equally well for elliptic systems). In (1.15) we have used $\Phi_{\varepsilon,k}^*$ to denote the Dirichlet correctors for $\mathcal{L}_{\varepsilon}^*$, the adjoint of $\mathcal{L}_{\varepsilon}$. Also, $h(y) = (\hat{a}_{ij}n_i(y)n_j(y))^{-1}$ and (\hat{a}_{ij}) is the (constant) coefficient matrix of \mathcal{L}_0 . The expansion (1.14) was used in [4] to identify the limit, as $\varepsilon \to 0$, of solutions to a problem of exact boundary controllability for the wave operator $\frac{\partial^2}{\partial t^2} + \mathcal{L}_{\varepsilon}$.

Our Theorem 1.1 gives a much more refined estimate of $R_{\varepsilon}(x,y)$ in (1.14) (under the stronger condition $\partial \Omega \in C^{2,\eta}$). Indeed, it follows from the estimate (1.8) that

$$|R_{\varepsilon}(x,y)| \le \frac{C \varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^d} \quad \text{for any } x \in \Omega, y \in \partial\Omega.$$
 (1.17)

Besides its applications to boundary control problems, estimate (1.17) may also be used to investigate the Dirichlet problem

$$\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = 0 \quad \text{in } \Omega \quad \text{and} \quad u_{\varepsilon}(x) = f(x, x/\varepsilon) \quad \text{on } \partial\Omega,$$
 (1.18)

where f(x, y) is 1-periodic in y. The Dirichlet problem (1.18) arises natually in the study of boundary layer phenomena and higher-order convergence in periodic homogenization (see e.g. [25, 1, 15, 14] and their references). Let v_{ε} be the solution to

$$\mathcal{L}_0(v_{\varepsilon}) = 0 \quad \text{in } \Omega \quad \text{and} \quad v_{\varepsilon} = f(x, x/\varepsilon)\omega_{\varepsilon}(x) \quad \text{on } \partial\Omega,$$
 (1.19)

where ω_{ε} is given by (1.15). It follows from the estimate (1.17) that

$$||u_{\varepsilon} - v_{\varepsilon}||_{L^{p}(\Omega)} = O((\varepsilon[\ln(1/\varepsilon)]^{2})^{1/p})$$
 for any $1 \le p < \infty$

(see Theorem 3.9). This effectively reduces the asymptotic problem (1.18) to the study of convergence properties of ω_{ε} on $\partial\Omega$, under various geometric conditions on Ω . This line of research will be developed in a future work.

We now describe the main ideas in the proof of Theorems 1.1 and 1.2. The basic tools in our approach are representation formulas by Green and Neumann functions, uniform estimates for Green functions established in [2],

$$|\nabla_x G_{\varepsilon}(x,y)| + |\nabla_y G_{\varepsilon}(x,y)| \le C|x-y|^{1-d}$$
 and $|\nabla_x \nabla_y G_{\varepsilon}(x,y)| \le C|x-y|^{-d}$, (1.20)

and the same estimates obtained in [20] for Neumann functions $N_{\varepsilon}(x,y)$. Let $D_r = D(x_0,r) = B(x_0,r) \cap \Omega$ and $\Delta_r = \Delta(x_0,r) = B(x_0,r) \cap \partial\Omega$ for some $x_0 \in \overline{\Omega}$ and $0 < r < r_0$. First, to establish (1.7), we will show that if p > d,

$$||u_{\varepsilon} - u_{0}||_{L^{\infty}(D_{1})} \le C \int_{D_{4}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon ||\nabla u_{0}||_{L^{\infty}(D_{4})} + C_{p} \varepsilon ||\nabla^{2} u_{0}||_{L^{p}(D_{4})}, \tag{1.21}$$

where $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0})$ in D_{4} and $u_{\varepsilon} = u_{0}$ on Δ_{4} . This is done by considering $w_{\varepsilon}(x) = u_{\varepsilon}(x) - u_{0}(x) - \varepsilon \chi(x/\varepsilon) \nabla u_{0}$ and using the Green representation formula and the observation that $\mathcal{L}_{\varepsilon}(w_{\varepsilon}) = \varepsilon \frac{\partial}{\partial x_{i}} \left(b_{ijk}(x/\varepsilon) \frac{\partial^{2} u_{0}}{\partial x_{j} \partial x_{k}} \right)$, where $b_{ijk}(y)$ is a bounded periodic function. Estimate (1.7) follows from (1.21) by a more or less standard argument (see subsection 3.1). Next, we show in subsection 3.2 that

$$\|\nabla u_{\varepsilon} - (\nabla \Phi_{\varepsilon})(\nabla u_{0})\|_{L^{\infty}(D_{r})} \leq \frac{C}{r^{d+1}} \int_{D_{4r}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon r^{-1} \|\nabla u_{0}\|_{L^{\infty}(D_{4r})} + C\varepsilon \ln[\varepsilon^{-1}r + 2] \|\nabla^{2}u_{0}\|_{L^{\infty}(D_{4r})} + C\varepsilon r^{\eta} \|\nabla^{2}u_{0}\|_{C^{0,\eta}(D_{4r})},$$
(1.22)

if $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0})$ in D_{4r} and $u_{\varepsilon} = u_{0}$ on Δ_{4r} . Estimate (1.8) follows easily from (1.22) by taking $u_{\varepsilon}(x) = G_{\varepsilon}(x, y_{0})$ and $u_{0}(x) = G_{0}(x, y_{0})$. By repeating the argument, estimate (1.22) also gives an asymptotic expansion for $\nabla_{x}\nabla_{y}G_{\varepsilon}(x, y)$ (see Theorem 3.11). To prove (1.22), we let

$$w_{\varepsilon} = u_{\varepsilon}(x) - u_{0}(x) - \left\{\Phi_{\varepsilon,j}^{\beta} - P_{j}^{\beta}\right\} \frac{\partial u_{0}^{\beta}}{\partial x_{j}}$$
(1.23)

and represent w_{ε} in D_r , using the Green function in \widetilde{D} , where \widetilde{D} is a $C^{2,\eta}$ domain such that $D_{3r} \subset \widetilde{D} \subset D_{4r}$.

Although a bit more complicated, the proof of Theorem 1.2 follows the same line of argument as Theorem 1.1. In subsections 4.1 and 4.2 we establish boundary L^{∞} and Lipschitz estimates similar to (1.21) and (1.22) for u_{ε} and u_0 satisfying $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in D_{4r} and $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\partial u_0}{\partial \nu_0}$ on Δ_{4r} . The results rely on the uniform L^p and Neumann function estimates obtained in [23, 21] under the additional symmetry condition $A^* = A$.

The rest of the paper is organized as follows. Section 2 contains some basic formulas and estimates which are more or less known. The case of Dirichlet boundary conditions is treated in Section 3, while Section 4 is devoted to the case of Neumann boundary conditions. In Section 5 we prove two inequalities, which are used in subsection 4.3 and are of interest in their own right, for the Dirichlet-to-Neumann map Λ_0 .

2 Preliminaries

Let $\mathcal{L}_{\varepsilon} = -\text{div}(A(x/\varepsilon)\nabla)$ with $A(y) = \left(a_{ij}^{\alpha\beta}(y)\right)$ satisfying (1.2)-(1.3). Let $\chi(y) = \left(\chi_j^{\alpha\beta}(y)\right)$ denote the matrix of correctors for \mathcal{L}_1 in \mathbb{R}^d , where $\chi_j^{\beta}(y) = \left(\chi_j^{1\beta}(y), \dots, \chi_j^{m\beta}(y)\right)$ is defined by the following cell problem:

$$\begin{cases} \mathcal{L}_1(\chi_j^{\beta}) = -\mathcal{L}_1(P_j^{\beta}) & \text{in } \mathbb{R}^d, \\ \chi_j^{\beta} \text{ is periodic with respect to } \mathbb{Z}^d \text{ and } \int_Y \chi_j^{\beta} dy = 0, \end{cases}$$
 (2.1)

for each $1 \leq j \leq d$ and $1 \leq \beta \leq m$. Here $Y = [0,1)^d \simeq \mathbb{R}^d/\mathbb{Z}^d$ and $P_j^{\beta}(y) = y_j(0,\ldots,1,\ldots,0)$ with 1 in the β^{th} position. The homogenized operator is given by $\mathcal{L}_0 = -\text{div}(\hat{A}\nabla)$, where $\hat{A} = (\hat{a}_{ij}^{\alpha\beta})$ and

$$\hat{a}_{ij}^{\alpha\beta} = \int_{Y} \left[a_{ij}^{\alpha\beta} + a_{ik}^{\alpha\gamma} \frac{\partial}{\partial y_{k}} \left(\chi_{j}^{\gamma\beta} \right) \right] dy. \tag{2.2}$$

It is known that the constant matrix \hat{A} is positive definite with an ellipticity constant depending only on d, m and μ (see [6]).

Let

$$b_{ij}^{\alpha\beta}(y) = \hat{a}_{ij}^{\alpha\beta} - a_{ij}^{\alpha\beta}(y) - a_{ik}^{\alpha\gamma}(y) \frac{\partial}{\partial y_k} (\chi_j^{\gamma\beta}). \tag{2.3}$$

Since $\int_Y b_{ij}^{\alpha\beta} dy = 0$ and $\frac{\partial}{\partial y_i} (b_{ij}^{\alpha\beta}) = 0$ by (2.2) and (2.1), there exists $F_{kij}^{\alpha\beta} \in H^1(Y)$ such that

$$b_{ij}^{\alpha\beta} = \frac{\partial}{\partial u_k} \left\{ F_{kij}^{\alpha\beta} \right\} \quad \text{and} \quad F_{kij}^{\alpha\beta} = -F_{ikj}^{\alpha\beta}.$$
 (2.4)

Remark 2.1. To see (2.4), one solves $\Delta f_{ij}^{\alpha\beta} = b_{ij}^{\alpha\beta}$ in Y with $f_{ij}^{\alpha\beta} \in H^1(Y)$ and $\int_Y f_{ij}^{\alpha\beta} dy = 0$, and let

$$F_{kij}^{\alpha\beta}(y) = \frac{\partial f_{ij}^{\alpha\beta}}{\partial y_k} - \frac{\partial f_{kj}^{\alpha\beta}}{\partial y_i}$$

(see e.g. [20]). Note that if A(y) is Hölder continuous, then $\nabla \chi$ and hence $b_{ij}^{\alpha\beta}$ are Hölder continuous. It follows that ∇F is Hölder continuous. In particular, $\|\chi\|_{C^1(Y)} + \|F\|_{C^1(Y)}$ is bounded by a constant depending only on d, m, μ , λ and τ . In the case of the scalar equation (m=1) with bounded measurable coefficients, the corrector χ is Hölder continuous by the De Giorgi-Nash estimates. This, together with Cacciopoli's inequality and Hölder's inequality, implies that there exist t>0 and C>0, depending only on d and μ , such that

$$\int_{B(x,r)} |\nabla \chi| \, dy \leq C \, r^{d-1+t} \quad \text{ for } x \in Y \text{ and } 0 < r < 1.$$

In view of (2.3) we obtain

$$\int_{B(x,r)} |b_{ij}(y)| \, dy \le C \, r^{d-1+t} \quad \text{for } x \in Y \text{ and } 0 < r < 1.$$
 (2.5)

Since $\Delta f_{ij} = b_{ij}$ in Y and $\int_Y f_{ij} dy = 0$,

$$\|\nabla f_{ij}\|_{L^{\infty}(Y)} \le C\|\nabla f_{ij}\|_{L^{2}(Y)} + C \sup_{x \in Y} \int_{Y} \frac{|b_{ij}(y)|}{|x - y|^{d - 1}} \, dy \le C(d, \mu), \tag{2.6}$$

where we have used (2.5) to estimate the last integral in (2.6). It follows that $||F_{kij}||_{\infty} \le C(d,\mu)$.

The following proposition plays an important role in this paper. We mention that formula (2.8) with $V_{\varepsilon,j}^{\beta}(x) = P_j^{\beta}(x) + \varepsilon \chi_j^{\beta}(x/\varepsilon)$ is known and may be used to show that $\|u_{\varepsilon} - u_0\|_{L^2(\Omega)} \le C\varepsilon \|u_0\|_{H^2(\Omega)}$, where $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in Ω and $u_{\varepsilon} = u_0$ on $\partial\Omega$ (see e.g. [20]). The proof of our main results on the first-order derivatives of Green and Neumann functions will rely on (2.8) with the matrices of Dirichlet and Neumann correctors respectively in the place of the functions $V_{\varepsilon,j}^{\beta}$.

Proposition 2.2. Suppose that $u_{\varepsilon} \in H^1(\Omega), u_0 \in H^2(\Omega)$ and $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in Ω . Let

$$w_{\varepsilon}(x) = u_{\varepsilon}(x) - u_{0}(x) - \left\{ V_{\varepsilon,j}^{\beta}(x) - P_{j}^{\beta}(x) \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}}, \tag{2.7}$$

where $V_{\varepsilon,j}^{\beta} = (V_{\varepsilon,j}^{1\beta}, \dots, V_{\varepsilon,j}^{m\beta}) \in H^1(\Omega)$ and $\mathcal{L}_{\varepsilon}(V_{\varepsilon,j}^{\beta}) = 0$ in Ω for each $1 \leq j \leq d$ and $1 \leq \beta \leq m$. Then

$$(\mathcal{L}_{\varepsilon}(w_{\varepsilon}))^{\alpha} = \varepsilon \frac{\partial}{\partial x_{i}} \left\{ \left[F_{jik}^{\alpha\gamma}(x/\varepsilon) \right] \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{j} \partial x_{k}} \right\}
+ \frac{\partial}{\partial x_{i}} \left\{ a_{ij}^{\alpha\beta}(x/\varepsilon) \left[V_{\varepsilon,k}^{\beta\gamma}(x) - x_{k} \delta^{\beta\gamma} \right] \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{j} \partial x_{k}} \right\}
+ a_{ij}^{\alpha\beta}(x/\varepsilon) \frac{\partial}{\partial x_{j}} \left[V_{\varepsilon,k}^{\beta\gamma}(x) - x_{k} \delta^{\beta\gamma} - \varepsilon \chi_{k}^{\beta\gamma}(x/\varepsilon) \right] \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{i} \partial x_{k}},$$
(2.8)

where $\delta^{\beta\gamma} = 1$ if $\beta = \gamma$, and zero otherwise.

Proof. Note that

$$a_{ij}^{\alpha\beta}(x/\varepsilon)\frac{\partial w_{\varepsilon}^{\beta}}{\partial x_{j}} = a_{ij}^{\alpha\beta}\left(\frac{x}{\varepsilon}\right)\frac{\partial u_{\varepsilon}^{\beta}}{\partial x_{j}} - a_{ij}^{\alpha\beta}\left(\frac{x}{\varepsilon}\right)\frac{\partial u_{0}^{\beta}}{\partial x_{j}} - a_{ij}^{\alpha\beta}\left(\frac{x}{\varepsilon}\right)\frac{\partial}{\partial x_{j}}\left\{V_{\varepsilon,k}^{\beta\gamma} - x_{k}\delta^{\beta\gamma}\right\} \cdot \frac{\partial u_{0}^{\gamma}}{\partial x_{k}} - a_{ij}^{\alpha\beta}\left(\frac{x}{\varepsilon}\right)\left\{V_{\varepsilon,k}^{\beta\gamma} - x_{k}\delta^{\beta\gamma}\right\} \cdot \frac{\partial^{2}u_{0}^{\gamma}}{\partial x_{k}\partial x_{j}}.$$

This, together with $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0})$, gives

$$\begin{aligned}
\left\{ \mathcal{L}_{\varepsilon}(w_{\varepsilon}) \right\}^{\alpha} &= -\frac{\partial}{\partial x_{i}} \left\{ \left[\hat{a}_{ij}^{\alpha\beta} - a_{ij}^{\alpha\beta} \left(x/\varepsilon \right) \right] \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \right\} + \left\{ \mathcal{L}_{\varepsilon} (V_{\varepsilon,k}^{\gamma} - P_{k}^{\gamma}) \right\}^{\alpha} \cdot \frac{\partial u_{0}^{\gamma}}{\partial x_{k}} \\
&+ a_{ij}^{\alpha\beta} \left(x/\varepsilon \right) \frac{\partial}{\partial x_{j}} \left\{ V_{\varepsilon,k}^{\beta\gamma} - x_{k} \delta^{\beta\gamma} \right\} \cdot \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{i} \partial x_{k}} \\
&+ \frac{\partial}{\partial x_{i}} \left\{ a_{ij}^{\alpha\beta} \left(x/\varepsilon \right) \left[V_{\varepsilon,k}^{\beta\gamma} - x_{k} \delta^{\beta\gamma} \right] \cdot \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{k} \partial x_{j}} \right\}.
\end{aligned}$$

Since

$$\mathcal{L}_{\varepsilon} \big(V_{\varepsilon,k}^{\gamma} - P_k^{\gamma} \big) = -\mathcal{L}_{\varepsilon} \big(P_k^{\gamma} \big) = \mathcal{L}_{\varepsilon} \big\{ \varepsilon \chi_k^{\gamma} (x/\varepsilon) \big\},$$

we obtain

$$\left\{ \mathcal{L}_{\varepsilon}(w_{\varepsilon}) \right\}^{\alpha} = -\frac{\partial}{\partial x_{i}} \left\{ b_{ij}^{\alpha\beta} \left(x/\varepsilon \right) \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \right\}
+ a_{ij}^{\alpha\beta} \left(x/\varepsilon \right) \frac{\partial}{\partial x_{j}} \left\{ V_{\varepsilon,k}^{\beta\gamma}(x) - x_{k} \delta^{\beta\gamma} - \frac{\partial \chi_{k}^{\beta\gamma}}{\partial x_{j}} \left(x/\varepsilon \right) \right\} \cdot \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{i} \partial x_{k}}
+ \frac{\partial}{\partial x_{i}} \left\{ a_{ij}^{\alpha\beta} \left(x/\varepsilon \right) \left[V_{\varepsilon,k}^{\beta\gamma} - x_{k} \delta^{\beta\gamma} \right] \cdot \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{k} \partial x_{j}} \right\},$$
(2.9)

where $b_{ij}^{\alpha\beta}(y)$ is defined by (2.3). In view of (2.4), we may re-write the first term in the right hand side of (2.9) as

$$-\frac{\partial}{\partial x_i} \left\{ \frac{\partial}{\partial x_k} \left[\varepsilon F_{kij}^{\alpha\beta} \left(x/\varepsilon \right) \right] \cdot \frac{\partial u_0^{\beta}}{\partial x_j} \right\} = \varepsilon \frac{\partial}{\partial x_i} \left\{ F_{kij}^{\alpha\beta} \left(x/\varepsilon \right) \cdot \frac{\partial^2 u_0^{\beta}}{\partial x_k \partial x_j} \right\}.$$

The formula (2.8) now follows.

The next proposition will be used to handle the Neumann boundary condition (cf. [21]).

Proposition 2.3. Let w_{ε} be given by (2.7). Suppose that $\frac{\partial}{\partial \nu_{\varepsilon}} \{V_{\varepsilon,k}^{\beta}\} = \frac{\partial}{\partial \nu_{0}} \{P_{k}^{\beta}\}$. Then

$$\left(\frac{\partial w_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} = \left(\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} - \left(\frac{\partial u_{0}}{\partial \nu_{0}}\right)^{\alpha} - n_{i} a_{ij}^{\alpha\beta}(x/\varepsilon) \left\{V_{\varepsilon,k}^{\beta\gamma} - P_{k}^{\beta\gamma}\right\} \cdot \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{k} \partial x_{j}}.$$
(2.10)

Proof. Note that

$$\left(\frac{\partial w_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} = \left(\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} - \left(\frac{\partial u_{0}}{\partial \nu_{\varepsilon}}\right)^{\alpha} - n_{i} a_{ij}^{\alpha\beta}(x/\varepsilon) \frac{\partial}{\partial x_{j}} \left[V_{\varepsilon,k}^{\beta\gamma} - P_{k}^{\beta\gamma}\right] \cdot \frac{\partial u_{0}^{\gamma}}{\partial x_{k}} - n_{i} a_{ij}^{\alpha\beta}(x/\varepsilon) \left[V_{\varepsilon,k}^{\beta\gamma} - P_{k}^{\beta\gamma}\right] \cdot \frac{\partial^{2} u_{0}^{\gamma}}{\partial x_{i} \partial x_{k}}.$$
(2.11)

Since
$$\frac{\partial}{\partial \nu_{\varepsilon}} \left\{ V_{\varepsilon,k}^{\beta} \right\} = \frac{\partial}{\partial \nu_{0}} \left\{ P_{k}^{\beta} \right\}$$
, the third term in the right hand side of (2.11) equals $-\left(\frac{\partial u_{0}}{\partial \nu_{0}} \right)^{\alpha} + \left(\frac{\partial u_{0}}{\partial \nu_{\varepsilon}} \right)^{\alpha}$. This gives (2.10).

The following two propositions provide the properties of the Dirichlet and Neumann correctors needed in this paper.

Proposition 2.4. Let $\mathcal{L}_{\varepsilon} = -\text{div}(A(x/\varepsilon)\nabla)$ with A(y) satisfying (1.2), (1.3) and (1.4). Let $(\Phi_{\varepsilon,i}^{\beta})$ denote the matrix of Dirichlet correctors for $\mathcal{L}_{\varepsilon}$ in a $C^{1,\eta}$ domain Ω . Then

$$|\nabla \Phi_j^{\beta}(x)| \le C, \qquad |\Phi_{\varepsilon,j}^{\beta}(x) - P_j^{\beta}(x)| \le C\varepsilon$$
 (2.12)

and

$$\left| \nabla \left\{ \Phi_{\varepsilon,j}^{\beta}(x) - P_{j}^{\beta}(x) - \varepsilon \chi_{j}^{\beta}(x/\varepsilon) \right\} \right| \leq C \min \left\{ 1, \frac{\varepsilon}{\delta(x)} \right\}$$
 (2.13)

for any $x \in \Omega$, where $\delta(x) = \operatorname{dist}(x, \partial\Omega)$ and C depends only d, m, μ, λ, τ and Ω .

Proof. The first estimate in (2.12) follows from the Lipschitz estimate in [2]. To see the second estimate, let $u_{\varepsilon}(x) = \Phi_{\varepsilon,j}^{\beta}(x) - P_{j}^{\beta}(x) - \varepsilon \chi_{j}^{\beta}(x/\varepsilon)$. Then $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = 0$ in Ω and $u_{\varepsilon}(x) = -\varepsilon \chi_{j}^{\beta}(x/\varepsilon)$ for $x \in \partial \Omega$. It again follows from [2] that $\|u_{\varepsilon}\|_{L^{\infty}(\Omega)} \leq C\|u_{\varepsilon}\|_{L^{\infty}(\partial \Omega)} \leq C\varepsilon$. Hence, $\|\Phi_{\varepsilon,j}^{\beta} - P_{j}^{\beta}\|_{L^{\infty}(\Omega)} \leq C\varepsilon + \|\varepsilon \chi_{j}^{\beta}\|_{\infty} \leq C\varepsilon$. Finally, note that $\|\nabla u_{\varepsilon}\|_{L^{\infty}(\Omega)} \leq C$. Also, by the interior estimate in [2] and $\|u_{\varepsilon}\|_{L^{\infty}(\Omega)} \leq C\varepsilon$, one obtains $|\nabla u_{\varepsilon}(x)| \leq C\varepsilon[\delta(x)]^{-1}$. This gives the estimate (2.13).

Proposition 2.5. Let $\mathcal{L}_{\varepsilon} = -\text{div}(A(x/\varepsilon)\nabla)$ with A(y) satisfying (1.2), (1.3), (1.4) and the symmetry condition $A^* = A$. Let $(\Psi_{\varepsilon,j}^{\beta})$ denote the matrix of Neumann correctors for $\mathcal{L}_{\varepsilon}$ in a $C^{1,\eta}$ domain Ω . Suppose $\Psi_{\varepsilon,j}^{\beta}(x_0) = P_j^{\beta}(x_0)$ for some $x_0 \in \Omega$. Then

$$|\nabla \Psi_j^{\beta}(x)| \le C, \qquad |\Psi_{\varepsilon,j}^{\beta}(x) - P_j^{\beta}(x)| \le C\varepsilon \ln[\varepsilon^{-1}M + 2]$$
 (2.14)

and

$$\left|\nabla\left\{\Psi_{\varepsilon,j}^{\beta}(x) - P_{j}^{\beta}(x) - \varepsilon\chi_{j}^{\beta}(x/\varepsilon)\right\}\right| \leq C \min\left\{1, \frac{\varepsilon}{\delta(x)}\right\}$$
 (2.15)

for any $x \in \Omega$, where $M = \operatorname{diam}(\Omega)$ and C depends only d, m, μ, λ, τ and Ω .

Proof. The estimate (2.15) as well as the first estimate in (2.14) was proved in [21]. To prove the second estimate in (2.14), we let $H_{\varepsilon,j}^{\beta} = \Psi_{\varepsilon,j}^{\beta}(x) - P_{j}^{\beta}(x) - \varepsilon \chi_{j}^{\beta}(x/\varepsilon)$. Since $|\nabla H_{\varepsilon,j}^{\beta}(x)| \leq C \min\left(1, \varepsilon[\delta(x)]^{-1}\right)$, by the Fundamental Theorem of Calculus, we may deduce that $|H_{\varepsilon,j}^{\beta}(x) - H_{\varepsilon,j}^{\beta}(x)| \leq C\varepsilon \ln[\varepsilon^{-1}M + 2]$ for any $x, y \in \Omega$. Since $|H_{\varepsilon,j}^{\beta}(x_0)| = \varepsilon|\chi_{j}^{\beta}(x_0)|$, we obtain $|H_{\varepsilon,j}^{\beta}(x)| \leq C\varepsilon \ln[\varepsilon^{-1}M + 2]$ for any $x \in \Omega$. This gives the desired estimate. \square

3 Asymptotic behavior of Green functions

The goal of this section is to prove Theorem 1.1. We also establish several convergence theorems for solutions with Dirichlet boundary conditions.

Let $\mathcal{L}_{\varepsilon} = -\text{div}(A(x/\varepsilon)\nabla)$ with A(y) satisfying conditions (1.2), (1.3) and (1.4). Let $G_{\varepsilon}(x,y)$ denote the matrix of Green's functions for $\mathcal{L}_{\varepsilon}$ in a bounded domain Ω . It follows from [2] that if Ω is $C^{1,\eta}$ for some $\eta \in (0,1)$,

$$\begin{cases}
|G_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^{d-2}} \min\left\{1, \frac{\delta(x)}{|x-y|}, \frac{\delta(y)}{|x-y|}, \frac{\delta(x)\delta(y)}{|x-y|^2}\right\}, \\
|\nabla_x G_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^{d-1}} \min\left\{1, \frac{\delta(y)}{|x-y|}\right\}, \\
|\nabla_y G_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^{d-1}} \min\left\{1, \frac{\delta(x)}{|x-y|}\right\}, \\
|\nabla_x \nabla_y G_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^d}
\end{cases} (3.1)$$

for any $x, y \in \Omega$, where $\delta(x) = \operatorname{dist}(x, \partial\Omega)$. These estimates, which are well known for second-order elliptic operators with constant coefficients, play an essential role in our approach to Theorem 1.1.

3.1 L^{∞} estimates

In this subsection we give the proof of the estimate (1.7). As a corollary of (1.7), we also establish an $O(\varepsilon)$ estimate for $||u_{\varepsilon} - u_0||_{L^p(\Omega)}$ for any p > 1 (see Theorem 3.4). Throughout

the subsection we will assume that A(y) satisfies conditions (1.2)-(1.3) and in the case m > 1, A(y) is Hölder continuous. Let

$$D_r = D_r(x_0, r) = B(x_0, r) \cap \Omega$$
 and $\Delta_r = \Delta(x_0, r) = B(x_0, r) \cap \partial\Omega$

for some $x_0 \in \overline{\Omega}$ and $0 < r < r_0$.

Lemma 3.1. Assume that Ω is Lipschitz if m=1, and $C^{1,\eta}$ for some $\eta \in (0,1)$ if m>1. Then

$$||u_{\varepsilon}||_{L^{\infty}(D_r)} \le C||f||_{L^{\infty}(\Delta_{3r})} + \frac{C}{r^d} \int_{D_{2r}} |u_{\varepsilon}| \, dx, \tag{3.2}$$

where $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = 0$ in D_{3r} and $u_{\varepsilon} = f$ on Δ_{3r} .

Proof. By rescaling we may assume that r=1. The estimate is well known in the case m=1 and follows from the maximum principle and De Giorgi -Nash estimate. If m>1 and f=0, estimate (3.2) follows directly from [2, Lemma 12]. To treat the general case, consider $\mathcal{L}_{\varepsilon}(v_{\varepsilon})=0$ in \widetilde{D} with the Dirichlet data $v_{\varepsilon}=f$ on $\partial \widetilde{D} \cap \partial \Omega$ and $v_{\varepsilon}=0$ on $\partial \widetilde{D} \setminus \partial \Omega$, where \widetilde{D} is a $C^{1,\eta}$ domain such that $D_2 \subset \widetilde{D} \subset D_3$. Note that $\|v_{\varepsilon}\|_{L^{\infty}(D_2)} \leq C\|f\|_{L^{\infty}(\Delta_3)}$ by [2, Theorem 3], and $u_{\varepsilon}-v_{\varepsilon}$ may be handled by [2, Lemma 12], as before.

The next lemma provides a boundary L^{∞} estimate.

Lemma 3.2. Assume that Ω satisfies the same assumption as in Lemma 3.1. Let $u_{\varepsilon} \in H^1(D_{4r})$ and $u_0 \in W^{2,p}(D_{4r})$ for some d . Suppose that

$$\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0})$$
 in D_{4r} and $u_{\varepsilon} = u_{0}$ on Δ_{4r} .

Then,

$$||u_{\varepsilon} - u_{0}||_{L^{\infty}(D_{r})} \leq \frac{C}{r^{d}} \int_{D_{4r}} |u_{\varepsilon} - u_{0}| dx + C\varepsilon ||\nabla u_{0}||_{L^{\infty}(D_{4r})} + C_{p} \varepsilon r^{1 - \frac{d}{p}} ||\nabla^{2} u_{0}||_{L^{p}(D_{4r})}.$$
(3.3)

Proof. Note that if $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$, then $\mathcal{L}_{\varepsilon/r}(v) = F_1$, where $v(x) = r^{-2}u_{\varepsilon}(rx)$ and $F_1(x) = F(rx)$. Thus, by rescaling, it suffices to consider the case r = 1. To this end we choose a domain \widetilde{D} , which is $C^{1,\eta}$ for m > 1 and Lipschitz for m = 1, such that $D_3 \subset \widetilde{D} \subset D_4$. Consider

$$w_{\varepsilon} = u_{\varepsilon} - u_0 - \varepsilon \chi_j^{\beta} (x/\varepsilon) \frac{\partial u_0^{\beta}}{\partial x_j} = w_{\varepsilon}^{(1)} + w_{\varepsilon}^{(2)} \quad \text{in } \widetilde{D},$$

where

$$\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(1)}) = \mathcal{L}_{\varepsilon}(w_{\varepsilon}) \quad \text{in } \widetilde{D} \quad \text{and} \quad w_{\varepsilon}^{(1)} \in H_0^1(\widetilde{D})$$
 (3.4)

and

$$\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(2)}) = 0 \quad \text{in } \widetilde{D} \quad \text{and} \quad w_{\varepsilon}^{(2)} = w_{\varepsilon} \quad \text{on } \partial \widetilde{D}.$$
 (3.5)

Since $w_{\varepsilon}^{(2)} = w_{\varepsilon} = -\varepsilon \chi(x/\varepsilon) \nabla u_0$ on Δ_3 and $\|\chi\|_{\infty} \leq C$, it follows from Lemma 3.1 that

$$||w_{\varepsilon}^{(2)}||_{L^{\infty}(D_{1})} \leq C\varepsilon ||\nabla u_{0}||_{L^{\infty}(\Delta_{3})} + C \int_{D_{3}} |w_{\varepsilon}^{(2)}| dx$$

$$\leq C\varepsilon ||\nabla u_{0}||_{L^{\infty}(\Delta_{3})} + C \int_{D_{3}} |w_{\varepsilon}| dx + C \int_{D_{3}} |w_{\varepsilon}^{(1)}| dx$$

$$\leq C \int_{D_{3}} |u_{\varepsilon} - u_{0}| dx + C\varepsilon ||\nabla u_{0}||_{L^{\infty}(D_{3})} + C ||w_{\varepsilon}^{(1)}||_{L^{\infty}(D_{3})}.$$

This gives

$$||u_{\varepsilon} - u_{0}||_{L^{\infty}(D_{1})} \le C \int_{D_{3}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon ||\nabla u_{0}||_{L^{\infty}(D_{3})} + C||w_{\varepsilon}^{(1)}||_{L^{\infty}(D_{3})}. \tag{3.6}$$

To estimate $w_{\varepsilon}^{(1)}$ in D_3 , we use the Green function representation

$$w_{\varepsilon}^{(1)}(x) = \int_{\widetilde{D}} \widetilde{G}_{\varepsilon}(x, y) \mathcal{L}_{\varepsilon}(w_{\varepsilon})(y) dy,$$

where $\widetilde{G}_{\varepsilon}(x,y)$ denotes the matrix of Green functions for $\mathcal{L}_{\varepsilon}$ in \widetilde{D} . Using (2.8) with $V_{\varepsilon,j}^{\beta} = P_{j}^{\beta}(x) + \varepsilon \chi_{j}^{\beta}(x/\varepsilon)$, we obtain

$$w_{\varepsilon}^{(1)}(x) = -\varepsilon \int_{\widetilde{D}} \frac{\partial}{\partial y_i} \left\{ \widetilde{G}_{\varepsilon}(x, y) \right\} \cdot \left[F_{jik} \left(y/\varepsilon \right) + a_{ij} \left(y/\varepsilon \right) \chi_k \left(y/\varepsilon \right) \right] \cdot \frac{\partial^2 u_0}{\partial y_i \partial y_k} \, dy,$$

where we have suppressed the superscripts for notational simplicity. Note that by Remark 2.1, $||F_{jik}||_{\infty} \leq C$. It follows that

$$|w_{\varepsilon}^{(1)}(x)| \leq C\varepsilon \int_{\widetilde{D}} |\nabla_{y}\widetilde{G}_{\varepsilon}(x,y)| |\nabla^{2}u_{0}(y)| dy$$

$$\leq C\varepsilon \|\nabla^{2}u_{0}\|_{L^{p}(D_{4})} \left\{ \int_{\widetilde{D}} |\nabla_{y}\widetilde{G}_{\varepsilon}(x,y)|^{p'} dy \right\}^{1/p'}$$

$$\leq C_{p} \varepsilon \|\nabla^{2}u_{0}\|_{L^{p}(D_{4})}$$

$$(3.7)$$

if p > d, where we have used Hölder's inequality and $\|\nabla_y \widetilde{G}_{\varepsilon}(x, \cdot)\|_{L^{p'}(\widetilde{D})} \leq C_p$. This, together with (3.6), completes the proof.

We are now ready to prove the first estimate in Theorem 1.1. Note that in the scalar case m = 1, no smoothness condition on A(y) is needed in the following theorem.

Theorem 3.3. Suppose that A(y) satisfies conditions (1.2)-(1.3). If m > 1, we also assume that A(y) is Hölder continuous. Let Ω be a bounded $C^{1,1}$ domain. Then

$$|G_{\varepsilon}(x,y) - G_0(x,y)| \le \frac{C\varepsilon}{|x-y|^{d-1}} \quad \text{for any } x, y \in \Omega,$$
 (3.8)

where C depends only on d, m, μ , Ω as well as λ and τ (if m > 1).

Proof. Under the assumptions on A and Ω , the estimates $|G_{\varepsilon}(x,y)| \leq C|x-y|^{2-d}$ and $|\nabla_x G_0(x,y)| \leq C|x-y|^{1-d}$ hold for any $x,y \in \Omega$ and $\varepsilon \geq 0$. We now fix $x_0,y_0 \in \Omega$ and $r = |x_0 - y_0|/8$. Let $f \in C_0^{\infty}(D(y_0,r))$,

$$u_{\varepsilon}(x) = \int_{\Omega} G_{\varepsilon}(x, y) f(y) dy$$
 and $u_{0}(x) = \int_{\Omega} G_{0}(x, y) f(y) dy$.

Then $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0}) = f$ in Ω and $u_{\varepsilon} = u_{0} = 0$ on $\partial\Omega$. Also, note that since Ω is $C^{1,1}$,

$$\begin{cases}
\|\nabla^2 u_0\|_{L^p(\Omega)} \le C_p \|f\|_{L^p(D(y_0,r))} & \text{for any } 1 d
\end{cases}$$
(3.9)

(see e. g. [16]). Next, let

$$w_{\varepsilon} = u_{\varepsilon} - u_0 - \varepsilon \chi_j^{\beta} (x/\varepsilon) \frac{\partial u_0^{\beta}}{\partial x_j} = \theta_{\varepsilon}(x) + z_{\varepsilon}(x),$$

where $\theta_{\varepsilon} \in H_0^1(\Omega)$ and $\mathcal{L}_{\varepsilon}(\theta_{\varepsilon}) = \mathcal{L}_{\varepsilon}(w_{\varepsilon})$ in Ω . Observe that by (2.8) with $V_{j,\varepsilon}^{\beta} = P_j^{\beta}(x) + \varepsilon \chi_j^{\beta}(x/\varepsilon)$ and Remark 2.1, $\|\theta_{\varepsilon}\|_{H_0^1(\Omega)} \leq C \varepsilon \|\nabla^2 u_0\|_{L^2(\Omega)} \leq C \varepsilon \|f\|_{L^2(D(y_0,r_0))}$. By Hölder's and Sobolev's inequalities, this implies that

$$\|\theta_{\varepsilon}\|_{L^{2}(D(x_{0},r))} \leq C r \|\theta_{\varepsilon}\|_{L^{q}(\Omega)} \leq C \varepsilon r \|f\|_{L^{2}(D(y_{0},r))} \leq C \varepsilon r^{1+\frac{d}{2}-\frac{d}{p}} \|f\|_{L^{p}(D(y_{0},r))},$$
(3.10)

where $q = \frac{2d}{d-2}$ and p > d. Also, note that since $\mathcal{L}_{\varepsilon}(z_{\varepsilon}) = 0$ in Ω and $z_{\varepsilon} = w_{\varepsilon}$ on $\partial\Omega$,

$$||z_{\varepsilon}||_{L^{\infty}(\Omega)} \le C||z_{\varepsilon}||_{L^{\infty}(\partial\Omega)} \le C \varepsilon ||\nabla u_{0}||_{L^{\infty}(\partial\Omega)}.$$
(3.11)

In view of (3.9)-(3.11), we obtain

$$||u_{\varepsilon} - u_{0}||_{L^{2}(D(x_{0},r))} \leq ||\theta_{\varepsilon}||_{L^{2}(D(x_{0},r))} + ||z_{\varepsilon}||_{L^{2}(D(x_{0},r))} + C \varepsilon r^{\frac{d}{2}} ||\nabla u_{0}||_{L^{\infty}(\Omega)}$$

$$\leq ||\theta_{\varepsilon}||_{L^{2}(D(x_{0},r))} + C \varepsilon r^{\frac{d}{2}} ||\nabla u_{0}||_{L^{\infty}(\Omega)}$$

$$\leq C \varepsilon r^{1 + \frac{d}{2} - \frac{d}{p}} ||f||_{L^{p}(D(y_{0},r))},$$
(3.12)

where p > d. This, together with Lemma 3.2 and (3.9), gives

$$|u_{\varepsilon}(x_0) - u_0(x_0)| \le C_p \, \varepsilon r^{1 - \frac{d}{p}} ||f||_{L^p(D(y_0, r))}.$$

It then follows by duality that

$$\left\{ \int_{D(y_0,r)} |G_{\varepsilon}(x_0,y) - G_0(x_0,y)|^{p'} dy \right\}^{1/p'} \le C_p \varepsilon r^{1-\frac{d}{p}} \quad \text{for any } p > d.$$

Finally, since $\mathcal{L}_{\varepsilon}^*(G_{\varepsilon}(x_0,\cdot)) = \mathcal{L}_0^*(G_0(x_0,\cdot)) = 0$ in $D(y_0,r)$, we may invoke Lemma 3.2 again to conclude that

$$|G_{\varepsilon}(x_{0}, y_{0}) - G_{0}(x_{0}, y_{0})| \leq \frac{C}{r^{d}} \int_{D(y_{0}, r)} |G_{\varepsilon}(x_{0}, y) - G_{0}(x_{0}, y)| dy$$

$$+ C \varepsilon ||\nabla_{y} G_{0}(x_{0}, \cdot)||_{L^{\infty}(D(y_{0}, r))}$$

$$+ C_{p} \varepsilon r^{1 - \frac{d}{p}} ||\nabla_{y}^{2} G_{0}(x_{0}, \cdot)||_{L^{p}(D(y_{0}, r_{0}))}$$

$$\leq C \varepsilon r^{1 - d},$$

where we also used

$$\left\{ \frac{1}{r^d} \int_{D(y_0,r)} |\nabla_y^2 G_0(x_0,y)|^p \, dy \right\}^{1/p} \le C_p \, r^{-2} \|G_0(x_0,\cdot)\|_{L^{\infty}(D(y_0,2r))} \le C_p \, r^{-d},$$

obtained by the boundary $W^{2,p}$ estimates on $C^{1,1}$ domains [16]. This completes the proof. \square

As a corollary of estimate (3.8), we obtain an $O(\varepsilon)$ estimate for $||u_{\varepsilon} - u_0||_{L^p(\Omega)}$ for any p > 1. In particular, we recover the estimate $||u_{\varepsilon} - u_0||_{L^2(\Omega)} \le C\varepsilon ||F||_{L^2(\Omega)}$, proved in [17] for scalar equations with bounded measurable coefficients satisfying (1.2)-(1.3). Also see [20] for estimates of $||u_{\varepsilon} - u_0||_{L^2(\Omega)}$ on Lipschitz domains.

Theorem 3.4. Suppose that A(y) and Ω satisfy the same conditions as in Theorem 3.3. For $F \in L^2(\Omega)$ and $\varepsilon \geq 0$, let $u_{\varepsilon} \in H^1_0(\Omega)$ be the solution of $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω . Then the estimate

$$||u_{\varepsilon} - u_0||_{L^q(\Omega)} \le C\varepsilon ||F||_{L^p(\Omega)} \tag{3.13}$$

holds if $1 and <math>\frac{1}{q} = \frac{1}{p} - \frac{1}{d}$, or p > d and $q = \infty$. Moreover,

$$||u_{\varepsilon} - u_0||_{L^{\infty}(\Omega)} \le C\varepsilon \left[\ln\left(\varepsilon^{-1}M + 2\right)\right]^{1 - \frac{1}{d}} ||F||_{L^d(\Omega)},\tag{3.14}$$

where $M = \operatorname{diam}(\Omega)$.

Proof. It follows from the Green function representation and Theorem 3.3 that

$$|u_{\varepsilon}(x) - u_0(x)| \le C \varepsilon \int_{\Omega} \frac{|F(y)|}{|x - y|^{d-1}} dy$$
, for any $x \in \Omega$.

This leads to (3.13) for $1 and <math>\frac{1}{q} = \frac{1}{p} - \frac{1}{d}$ by the well known estimates for fractional integrals. The case of p > d and $q = \infty$ follows directly from Hölder's inequality. To see (3.14), we bound $|G_{\varepsilon}(x,y) - G_0(x,y)|$ by $C|x-y|^{2-d}$ if $|x-y| < \varepsilon$, and by $C\varepsilon|x-y|^{1-d}$ if $|x-y| \ge \varepsilon$. By Hölder's inequality, this gives

$$|u_{\varepsilon}(x) - u_{0}(x)| \leq C \int_{D(x,\varepsilon)} \frac{|F(y)|}{|x - y|^{d-2}} dy + C\varepsilon \int_{\Omega \setminus D(x,\varepsilon)} \frac{|F(y)|}{|x - y|^{d-1}} dy$$

$$\leq C\varepsilon ||F||_{L^{d}(\Omega)} + C\varepsilon \left[\ln\left(\varepsilon^{-1}M + 2\right)\right]^{1 - \frac{1}{d}} ||F||_{L^{d}(\Omega)}$$

$$\leq C\varepsilon \left[\ln\left(\varepsilon^{-1}M + 2\right)\right]^{1 - \frac{1}{d}} ||F||_{L^{d}(\Omega)},$$

which completes the proof.

3.2 Lipschitz estimates

In this subsection we give the proof of (1.8). As a corollary of (1.8), we also obtain an $O(\varepsilon)$ estimate for $u_{\varepsilon} - u_0 - \{\Phi_{\varepsilon,j} - P_j\}\frac{\partial u_0}{\partial x_j}$ in $W_0^{1,p}(\Omega)$ for any 1 .

Recall that $D(r) = D(x_0, r) = B(x_0, r) \cap \Omega$ and $\Delta_r = \Delta(x_0, r) = B(x_0, r) \cap \partial \Omega$, where $x_0 \in \overline{\Omega}$ and $0 < r < r_0$. Throughout this subsection we will assume that Ω is a bounded $C^{2,\eta}$ domain for some $\eta \in (0, 1)$ and A = A(y) satisfies conditions (1.2), (1.3) and (1.4).

Lemma 3.5. Suppose that $u_{\varepsilon} \in H^1(D_{4r})$, $u_0 \in C^{2,\rho}(D_{4r})$ and $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in D_{4r} , where $0 < \rho < \eta$. Also assume that $u_{\varepsilon} = u_0$ on Δ_{4r} . Then, if $0 < \varepsilon < r$,

$$\|\frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon,j}^{\alpha\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \|_{L^{\infty}(D_{r})}$$

$$\leq \frac{C}{r^{d+1}} \int_{D_{4r}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon r^{-1} \|\nabla u_{0}\|_{L^{\infty}(D_{4r})}$$

$$+ C\varepsilon \ln \left[\varepsilon^{-1} r + 2 \right] \|\nabla^{2} u_{0}\|_{L^{\infty}(D_{4r})} + C\varepsilon r^{\rho} \|\nabla^{2} u_{0}\|_{C^{0,\rho}(D_{4r})},$$

$$(3.15)$$

where $\Phi_{\varepsilon} = (\Phi_{\varepsilon,j}^{\alpha\beta})$ denotes the matrix of Dirichlet correctors for $\mathcal{L}_{\varepsilon}$ in Ω .

Proof. We begin by choosing a $C^{2,\eta}$ domain \widetilde{D} such that $D_{3r} \subset \widetilde{D} \subset D_{4r}$. Let

$$w_{\varepsilon} = u_{\varepsilon} - u_0 - \left\{ \Phi_{\varepsilon,j}^{\beta} - P_j^{\beta} \right\} \cdot \frac{\partial u_0^{\beta}}{\partial x_j}.$$

Note that $w_{\varepsilon} = 0$ on Δ_{4r} . Write $w_{\varepsilon} = w_{\varepsilon}^{(1)} + w_{\varepsilon}^{(2)}$ in \widetilde{D} , where $w_{\varepsilon}^{(1)} \in H_0^1(\widetilde{D})$ and $\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(1)}) = \mathcal{L}_{\varepsilon}(w_{\varepsilon})$ in \widetilde{D} . Since $\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(2)}) = 0$ in D_{3r} and $w_{\varepsilon}^{(2)} = w_{\varepsilon} = 0$ on Δ_{3r} , it follows from the boundary Lipschitz estimate in [2, Lemma 20] that

$$\begin{split} \|\nabla w_{\varepsilon}^{(2)}\|_{L^{\infty}(D_{r})} &\leq \frac{C}{r^{d+1}} \int_{D_{2r}} |w_{\varepsilon}^{(2)}| \, dx \\ &\leq \frac{C}{r^{d+1}} \int_{D_{2r}} |w_{\varepsilon}| \, dx + Cr^{-1} \|w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{2r})} \\ &\leq \frac{C}{r^{d+1}} \int_{D_{2r}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon r^{-1} \|\nabla u_{0}\|_{L^{\infty}(D_{2r})} + Cr^{-1} \|w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{2r})}, \end{split}$$

where we have used the estimate $\|\Phi_{\varepsilon,j}^{\beta} - P_j^{\beta}\|_{\infty} \leq C\varepsilon$ in Proposition 2.4. This implies that

$$\|\nabla w_{\varepsilon}\|_{L^{\infty}(D_r)} \leq \frac{C}{r^{d+1}} \int_{D_{2r}} |u_{\varepsilon} - u_0| \, dx + C\varepsilon r^{-1} \|\nabla u_0\|_{L^{\infty}(D_{2r})} + C \|\nabla w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{2r})},$$

where we have used $||w_{\varepsilon}^{(1)}||_{L^{\infty}(D_{2r})} \leq Cr ||\nabla w_{\varepsilon}^{(1)}||_{L^{\infty}(D_{2r})}$. Thus,

$$\|\frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon,j}^{\alpha\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \|_{L^{\infty}(D_{r})}$$

$$\leq \frac{C}{r^{d+1}} \int_{D_{2r}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon r^{-1} \|\nabla u_{0}\|_{L^{\infty}(D_{2r})}$$

$$+ C\varepsilon \|\nabla^{2} u_{0}\|_{L^{\infty}(D_{2r})} + C \|\nabla w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{2r})}.$$

$$(3.16)$$

To estimate $\nabla w_{\varepsilon}^{(1)}$ on D_{2r} , we use the Green function representation

$$w_{\varepsilon}^{(1)}(x) = \int_{\widetilde{D}} \widetilde{G}_{\varepsilon}(x, y) \mathcal{L}_{\varepsilon}(w_{\varepsilon})(y) dy,$$

where $\widetilde{G}_{\varepsilon}(x,y)$ is the matrix of Green functions for $\mathcal{L}_{\varepsilon}$ in the $C^{2,\eta}$ domain \widetilde{D} . Let

$$f_i(x) = \varepsilon F_{kij}(x/\varepsilon) \frac{\partial^2 u_0}{\partial x_j \partial x_k} + a_{ij}(x/\varepsilon) \left[\Phi_{\varepsilon,k} - P_k \right] \cdot \frac{\partial^2 u_0}{\partial x_j \partial x_k},$$

where we have suppressed the superscripts for notational simplicity. In view of (2.8), we obtain

$$w_{\varepsilon}^{(1)}(x) = -\int_{\widetilde{D}} \frac{\partial}{\partial y_{i}} \left\{ \widetilde{G}_{\varepsilon}(x,y) \right\} \cdot \left\{ f_{i}(y) - f_{i}(x) \right\} dy$$
$$+ \int_{\widetilde{D}} \widetilde{G}_{\varepsilon}(x,y) a_{ij} \left(y/\varepsilon \right) \frac{\partial}{\partial y_{j}} \left[\Phi_{\varepsilon,k} - P_{k} - \varepsilon \chi_{k} \left(y/\varepsilon \right) \right] \cdot \frac{\partial^{2} u_{0}}{\partial y_{i} \partial y_{k}} dy.$$

It follows that

$$|\nabla w_{\varepsilon}^{(1)}(x)| \leq \int_{\widetilde{D}} |\nabla_{x} \nabla_{y} \widetilde{G}_{\varepsilon}(x, y)| |f(y) - f(x)| dy + C ||\nabla^{2} u_{0}||_{L^{\infty}(D_{4r})} \int_{\widetilde{D}} |\nabla_{x} \widetilde{G}_{\varepsilon}(x, y)| |\nabla_{y} [\Phi_{\varepsilon, j} - P_{j} - \varepsilon \chi_{j} (y/\varepsilon)]| dy.$$
(3.17)

To handle the first term in the right hand side of (3.17), we use $|\nabla_x \nabla_y \widetilde{G}_{\varepsilon}(x,y)| \leq C|x-y|^{-d}$ and the observation that

$$||f||_{L^{\infty}(D_{4r})} \le C \varepsilon ||\nabla^{2} u_{0}||_{L^{\infty}(D_{4r})},$$

$$|f(x) - f(y)| \le C|x - y|^{\rho} \left\{ \varepsilon^{1-\rho} ||\nabla^{2} u_{0}||_{L^{\infty}(D_{4r})} + \varepsilon ||\nabla^{2} u_{0}||_{C^{0,\rho}(D_{4r})} \right\}.$$

This yields that

$$\int_{\widetilde{D}} |\nabla_{x} \nabla_{y} \widetilde{G}_{\varepsilon}(x,y)| |f(y) - f(x)| dy$$

$$\leq C \varepsilon \|\nabla^{2} u_{0}\|_{L^{\infty}(D_{4r})} \int_{\widetilde{D} \setminus B(x,\varepsilon)} \frac{dy}{|x - y|^{d}}$$

$$+ C \left\{ \varepsilon^{1-\rho} \|\nabla^{2} u_{0}\|_{L^{\infty}(D_{4r})} + \varepsilon \|\nabla^{2} u_{0}\|_{C^{0,\rho}(D_{4r})} \right\} \int_{\widetilde{D} \cap B(x,\varepsilon)} \frac{dy}{|x - y|^{d-\rho}}$$

$$\leq C \varepsilon \ln[\varepsilon^{-1} r + 2] \|\nabla^{2} u_{0}\|_{L^{\infty}(D_{4r})} + C\varepsilon^{1+\rho} \|\nabla^{2} u_{0}\|_{C^{0,\rho}(D_{4r})}.$$

Finally, using the estimates $|\nabla_x \widetilde{G}_{\varepsilon}(x,y)| \leq C \operatorname{dist}(y,\partial\Omega)|x-y|^{-d}$ and $|\nabla_x \widetilde{G}_{\varepsilon}(x,y)| \leq C|x-y|^{1-d}$ as well as estimates in Proposition 2.4, we see that the second term in the right-hand side of (3.17) is bounded by

$$C\|\nabla^2 u_0\|_{L^{\infty}(D_{4r})} \left\{ \varepsilon \int_{\widetilde{D} \setminus B(x,\varepsilon)} \frac{dy}{|x-y|^d} + \int_{\widetilde{D} \cap B(x,\varepsilon)} \frac{dy}{|x-y|^{d-1}} \right\}$$

$$\leq C \varepsilon \ln[\varepsilon^{-1}r + 2] \|\nabla^2 u_0\|_{L^{\infty}(D_{4r})}.$$

Thus we have proved that

$$\|\nabla w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{3r})} \le C \varepsilon \ln[\varepsilon^{-1}r + 2] \|\nabla^{2}u_{0}\|_{L^{\infty}(D_{4r})} + C\varepsilon^{1+\rho} \|\nabla^{2}u_{0}\|_{C^{0,\rho}(D_{4r})}.$$

This, together with (3.16), completes the proof of (3.15).

We are now ready to give the proof of the estimate (1.8).

Theorem 3.6. Suppose that A(y) satisfies conditions (1.2), (1.3) and (1.4). Let Ω be a bounded $C^{2,\eta}$ domain for some $\eta \in (0,1)$. Then

$$\left| \frac{\partial}{\partial x_i} \left\{ G_{\varepsilon}^{\alpha \gamma}(x, y) \right\} - \frac{\partial}{\partial x_i} \left\{ \Phi_{\varepsilon, j}^{\alpha \beta}(x) \right\} \cdot \frac{\partial}{\partial x_j} \left\{ G_0^{\beta \gamma}(x, y) \right\} \right| \le \frac{C \varepsilon \ln[\varepsilon^{-1} |x - y| + 2]}{|x - y|^d}, \tag{3.18}$$

for any $x, y \in \Omega$, where C depends only on d, m, μ , λ , τ and Ω .

Proof. Fix $x_0, y_0 \in \Omega$ and $r = |x_0 - y_0|/8$. We may assume that $\varepsilon \leq r$, since the case $\varepsilon > r$ is trivial and follows from the size estimate of $|\nabla_x G_{\varepsilon}(x, y)|$, $|\nabla_x G_0(x, y)|$ and (2.12).

Let $u_{\varepsilon}(x) = G_{\varepsilon}(x, y_0)$ and $u_0(x) = G_0(x, y_0)$. Then $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0) = 0$ in $D_{4r} = D(x_0, 4r)$ and $u_{\varepsilon} = u_0 = 0$ on $\Delta_{4r} = \Delta(x_0, 4r)$. Note that by Theorem 3.3, we have $\|u_{\varepsilon} - u_0\|_{L^{\infty}(D_{4r})} \leq C \, \varepsilon r^{1-d}$. Also, since Ω is $C^{2,\eta}$, we have $\|\nabla u_0\|_{L^{\infty}(D_{4r})} \leq C r^{1-d}$, $\|\nabla^2 u_0\|_{L^{\infty}(D_{4r})} \leq C r^{-d-\rho}$. It then follows from Lemma 3.5 that

$$\left\| \frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon, j}^{\alpha \beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{i}} \right\|_{L^{\infty}(D_{r})} \leq C \varepsilon r^{-d} \ln[\varepsilon^{-1}r + 2].$$

This finishes the proof.

As a corollary of the estimate (3.18), we obtain an $O(\varepsilon)$ estimate (up to a logarithmic factor if $p \neq 2$) for $u_{\varepsilon} - u_0 - \{\Phi_{\varepsilon,j} - P_j\}\frac{\partial u_0}{\partial x_j}$ in $W_0^{1,p}(\Omega)$.

Theorem 3.7. Assume that A(y) and Ω satisfy the same assumptions as in Theorem 3.6. Let $1 . For <math>F \in L^p(\Omega)$ and $\varepsilon \ge 0$, let $u_{\varepsilon} \in W_0^{1,p}(\Omega)$ and $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω . Then

$$\|u_{\varepsilon} - u_0 - \left\{\Phi_{\varepsilon,j}^{\beta} - P_j^{\beta}\right\} \frac{\partial u_0^{\beta}}{\partial x_i} \|_{W_0^{1,p}(\Omega)} \le C_p \varepsilon \left\{\ln[\varepsilon^{-1}M + 2]\right\}^{4|\frac{1}{2} - \frac{1}{p}|} \|F\|_{L^p(\Omega)}, \tag{3.19}$$

where $M = \operatorname{diam}(\Omega)$ and C_p depends only on d, m, p, μ , λ , τ and Ω .

Proof. We will show that for any $1 \le p \le \infty$,

$$\left\| \frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon, j}^{\alpha \beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \right\|_{L^{p}(\Omega)} \le C \varepsilon \left\{ \ln[\varepsilon^{-1} M + 2] \right\}^{4\left|\frac{1}{2} - \frac{1}{p}\right|} \|F\|_{L^{p}(\Omega)}. \tag{3.20}$$

This, together with (2.12) and the estimate $\|\nabla^2 u_0\|_{L^p(\Omega)} \le C\|F\|_{L^p(\Omega)}$ for $1 , gives (3.19). To see (3.20), we use Theorem 3.6 as well as estimates on <math>|\nabla_x G_{\varepsilon}(x,y)|$ and $|\nabla \Phi_{\varepsilon}|$ to deduce that

$$\left| \frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon,j}^{\alpha\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \right| \leq C \int_{\Omega} K_{\varepsilon}(x,y) |f(y)| \, dy,$$

where

$$K_{\varepsilon}(x,y) = \begin{cases} \varepsilon |x-y|^{-d} \ln \left[\varepsilon^{-1} |x-y| + 2 \right], & \text{if } |x-y| \ge \varepsilon, \\ |x-y|^{1-d}, & \text{if } |x-y| < \varepsilon. \end{cases}$$
(3.21)

Note that

$$\sup_{x \in \Omega} \int_{\Omega} K_{\varepsilon}(x, y) \, dy + \sup_{y \in \Omega} \int_{\Omega} K_{\varepsilon}(x, y) \, dx \le C \varepsilon \big\{ \ln[\varepsilon^{-1} M + 2] \big\}^{2}.$$

This gives (3.20) in the case p = 1 or ∞ . Thus, by the M. Riesz interpolation theorem, it suffices to prove the estimate for p = 2.

Let $w_{\varepsilon} = u_{\varepsilon} - u_0 - \{\Phi_{\varepsilon,j}^{\beta} - P_j^{\beta}\}\frac{\partial u_0^{\beta}}{\partial x_j}$ and $\delta(x) = \operatorname{dist}(x,\partial\Omega)$. We may deduce from Propositions 2.2 and 2.4 that

$$\int_{\Omega} |\nabla w_{\varepsilon}|^{2} dx \leq C\varepsilon \int_{\Omega} |\nabla^{2} u_{0}| |\nabla w_{\varepsilon}| dx + C\varepsilon \int_{\Omega} [\delta(x)]^{-1} |\nabla^{2} u_{0}| |w_{\varepsilon}| dx
\leq C\varepsilon ||\nabla^{2} u_{0}||_{L^{2}(\Omega)} ||\nabla w_{\varepsilon}||_{L^{2}(\Omega)},$$
(3.22)

where, for the last inequality, we have used Hölder's inequality as well as Hardy's inequality $\|[\delta(x)]^{-1}w_{\varepsilon}\|_{L^{2}(\Omega)} \leq C\|\nabla w_{\varepsilon}\|_{L^{2}(\Omega)}$. The desired estimate follows easily from (3.22).

Let $G_{\varepsilon}^*(x,y) = (G_{\varepsilon}^{*\alpha\beta}(x,y))$ denote the matrix of Green's functions for $\mathcal{L}_{\varepsilon}^*$, the adjoint of $\mathcal{L}_{\varepsilon}$. By Theorem 3.6,

$$\left| \frac{\partial}{\partial x_i} \left\{ G_{\varepsilon}^{*\alpha\gamma}(x,y) \right\} - \frac{\partial}{\partial x_i} \left\{ \Phi_{\varepsilon,j}^{*\alpha\beta}(x) \right\} \cdot \frac{\partial}{\partial x_j} \left\{ G_0^{*\beta\gamma}(x,y) \right\} \right| \le \frac{C \varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^d}, \quad (3.23)$$

where Φ_{ε}^* denotes the matrix of Dirichlet correctors for $\mathcal{L}_{\varepsilon}^*$ in Ω . Since $G_{\varepsilon}^{*\alpha\beta}(x,y) = G_{\varepsilon}^{\beta\alpha}(y,x)$, we obtain

$$\left| \frac{\partial}{\partial y_i} \left\{ G_{\varepsilon}^{\gamma \alpha}(x, y) \right\} - \frac{\partial}{\partial y_i} \left\{ \Phi_{\varepsilon, j}^{*\alpha \beta}(y) \right\} \cdot \frac{\partial}{\partial y_i} \left\{ G_0^{\gamma \beta}(x, y) \right\} \right| \le \frac{C \varepsilon \ln[\varepsilon^{-1} |x - y| + 2]}{|x - y|^d}. \tag{3.24}$$

This leads to an asymptotic expansion of the Poisson kernel for $\mathcal{L}_{\varepsilon}$ on Ω .

Let $(h^{\alpha\beta}(y))$ denote the inverse matrix of $(n_i(y)n_j(y)\hat{a}_{ij}^{\alpha\beta})_{m\times m}$.

Theorem 3.8. Suppose that A(y) satisfies conditions (1.2), (1.3) and (1.4). Let $P_{\varepsilon}(x,y) = (P_{\varepsilon}^{\alpha\beta}(x,y))$ be the Poisson kernel for $\mathcal{L}_{\varepsilon}$ on a $C^{2,\eta}$ domain Ω . Then

$$P_{\varepsilon}^{\alpha\beta}(x,y) = P_0^{\alpha\gamma}(x,y)\omega_{\varepsilon}^{\gamma\beta}(y) + R_{\varepsilon}^{\alpha\beta}(x,y), \tag{3.25}$$

where

$$\omega_{\varepsilon}^{\gamma\beta}(y) = h^{\gamma\sigma}(y) \cdot \frac{\partial}{\partial n(y)} \left\{ \Phi_{\varepsilon,k}^{*\rho\sigma}(y) \right\} \cdot n_k(y) \cdot n_i(y) n_j(y) a_{ij}^{\rho\beta}(y/\varepsilon), \tag{3.26}$$

$$|R_{\varepsilon}^{\alpha\beta}(x,y)| \le \frac{C \varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^d} \quad \text{for any } x \in \Omega \text{ and } y \in \partial\Omega,$$
 (3.27)

and C depends only on d, m, μ , λ , τ and Ω .

Proof. Note that

$$P_{\varepsilon}^{\alpha\beta}(x,y) = -n_{i}(y)a_{ji}^{\gamma\beta}(y/\varepsilon)\frac{\partial}{\partial y_{j}}\left\{G_{\varepsilon}^{\alpha\gamma}(x,y)\right\}$$

$$= -\frac{\partial}{\partial n(y)}\left\{G_{\varepsilon}^{\alpha\gamma}(x,y)\right\} \cdot n_{i}(y)n_{j}(y)a_{ij}^{\gamma\beta}(y/\varepsilon),$$
(3.28)

since $G_{\varepsilon}(x,\cdot) = 0$ on $\partial\Omega$. By (3.24), we obtain

$$\left| P_{\varepsilon}^{\alpha\beta}(x,y) + \frac{\partial}{\partial n(y)} \left\{ G_0^{\alpha\sigma}(x,y) \right\} \cdot \frac{\partial}{\partial n(y)} \left\{ \Phi_{\varepsilon,k}^{*\gamma\sigma}(y) \right\} \cdot n_i(y) n_j(y) a_{ij}^{\gamma\beta}(y/\varepsilon) n_k(y) \right| \\
\leq \frac{C \varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^d}.$$
(3.29)

In view of (3.28) (with $\varepsilon = 0$), we have

$$P_0^{\alpha\beta}(x,y)h^{\beta\sigma}(y) = -\frac{\partial}{\partial n(y)} \{G_0^{\alpha\sigma}(x,y)\}.$$

This, together with (3.29), gives

$$|P_{\varepsilon}^{\alpha\beta}(x,y) - P_0^{\alpha\gamma}(x,y)\omega_{\varepsilon}^{\gamma\beta}(y)| \le \frac{C\,\varepsilon\ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^d},$$

for any $x \in \Omega$ and $y \in \partial\Omega$, where $\omega_{\varepsilon}(y)$ is defined by (3.26).

With the asymptotic expansion for the Poisson kernels at our disposal, we may approximate the solution of the L^p Dirichlet problem: $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = 0$ in Ω and $u_{\varepsilon} = f_{\varepsilon}$ on $\partial\Omega$ by the solution of the homogenized system with boundary data $\omega_{\varepsilon}f_{\varepsilon}$. As mentioned in the Introduction, the case $f_{\varepsilon}(x) = f(x, x/\varepsilon)$ with f(x, y) periodic in y is of particular interest.

Theorem 3.9. Assume that A(y) and Ω satisfy the same assumptions as in Theorem 3.8. Let $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = 0$ in Ω and $u_{\varepsilon} = f_{\varepsilon}$ on $\partial\Omega$. Then for any $1 \leq p < \infty$,

$$||u_{\varepsilon} - v_{\varepsilon}||_{L^{p}(\Omega)} \le C \left\{ \varepsilon \left(\ln[\varepsilon^{-1}M + 2] \right)^{2} \right\}^{1/p} ||f_{\varepsilon}||_{L^{p}(\partial\Omega)}, \tag{3.30}$$

where $\mathcal{L}_0(v_{\varepsilon}) = 0$ in Ω and $v_{\varepsilon} = \omega_{\varepsilon} f_{\varepsilon}$ on $\partial \Omega$, with ω_{ε} defined by (3.26).

Proof. Since

$$u_{\varepsilon}(x) = \int_{\partial\Omega} P_{\varepsilon}(x,y) f_{\varepsilon}(y) dy$$
 and $v_{\varepsilon}(x) = \int_{\partial\Omega} P_{0}(x,y) \omega_{\varepsilon}(y) f_{\varepsilon}(y) dy$,

it follows from Theorem 3.8 that

$$|u_{\varepsilon}(x) - v_{\varepsilon}(x)| \le \int_{\partial\Omega} |R_{\varepsilon}(x, y)| |f_{\varepsilon}(y)| dy.$$

Using

$$|R_{\varepsilon}(x,y)| \le C\{|\nabla_y G_{\varepsilon}(x,y)| + |\nabla_y G_0(x,y)|\} \le C\delta(x)|x-y|^{-d},$$

where $\delta(x) = \operatorname{dist}(x, \partial\Omega)$, we see that $\int_{\partial\Omega} |R_{\varepsilon}(x, y)| dy \leq C$ for any $x \in \Omega$. It then follows by Hölder's inequality that

$$|u_{\varepsilon}(x) - v_{\varepsilon}(x)|^p \le C \int_{\partial\Omega} |R_{\varepsilon}(x, y)| |f_{\varepsilon}(y)|^p \, dy. \tag{3.31}$$

Now, by Theorem 3.8 as well as the estimate $|R_{\varepsilon}(x,y)| \leq C|x-y|^{1-d}$, we obtain

$$\int_{\Omega} |R_{\varepsilon}(x,y)| \, dx \leq C \int_{\Omega \cap B(y,\varepsilon)} \frac{dx}{|x-y|^{d-1}} + C\varepsilon \int_{\Omega \setminus B(y,\varepsilon)} \frac{\ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^d} \, dx
\leq C\varepsilon \left\{ \ln[\varepsilon^{-1}M+2] \right\}^2,$$

for any $y \in \partial \Omega$. This, together with (3.31) and Fubini's Theorem, gives (3.30).

We end this section with another approximation result. Note that by (3.20),

$$\|\nabla(\mathcal{L}_{\varepsilon})^{-1}(F) - \nabla\Phi_{\varepsilon} \cdot \nabla(\mathcal{L}_{0})^{-1}(F)\|_{L^{p}(\Omega)} \le C\varepsilon \left\{ \ln[\varepsilon^{-1}M + 2] \right\}^{4\left|\frac{1}{2} - \frac{1}{p}\right|} \|F\|_{L^{p}(\Omega)}, \tag{3.32}$$

for $1 \le p \le \infty$. By duality this gives the following.

Theorem 3.10. Assume that A(y) and Ω satisfy the same assumptions as in Theorem 3.8. For $f = (f_i^{\alpha}) \in L^2(\Omega)$ and $\varepsilon \geq 0$, let $u_{\varepsilon} \in H_0^1(\Omega)$ and $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \operatorname{div}(f)$ in Ω . Then if $f \in L^p(\Omega)$ for some $1 \leq p \leq \infty$,

$$||u_{\varepsilon} - v_{\varepsilon}||_{L^{p}(\Omega)} \le C \varepsilon \left\{ \ln[\varepsilon^{-1}M + 2] \right\}^{4\left|\frac{1}{2} - \frac{1}{p}\right|} ||f||_{L^{p}(\Omega)}, \tag{3.33}$$

where $v_{\varepsilon} \in H_0^1(\Omega)$ and $\mathcal{L}_0(v_{\varepsilon}) = \operatorname{div}(F_{\varepsilon})$ with

$$F_{\varepsilon,i}^{\alpha}(x) = f_j^{\beta}(x) \frac{\partial}{\partial x_i} \left\{ \Phi_{\varepsilon,i}^{*\beta\alpha} \right\}.$$

3.3 An asymptotic expansion for $\nabla_x \nabla_y G_{\varepsilon}(x,y)$ and its applications

In this subsection we derive an asymptotic expansion for $\nabla_x \nabla_y G_{\varepsilon}(x,y)$. As its applications we obtain asymptotic expansions for $(\partial/\partial x_i)(\mathcal{L}_{\varepsilon})^{-1}(\partial/\partial x_j)$ and the Dirichlet-to-Neumann map associated with $\mathcal{L}_{\varepsilon}$.

Theorem 3.11. Suppose that A(y) satisfies conditions (1.2), (1.3) and (1.4). Let Ω be a $C^{3,\eta}$ domain for some $\eta > 0$. Then

$$\left| \frac{\partial^{2}}{\partial x_{i} \partial y_{j}} \left\{ G_{\varepsilon}^{\alpha \beta}(x, y) \right\} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon, k}^{\alpha \gamma}(x) \right\} \cdot \frac{\partial^{2}}{\partial x_{k} \partial y_{\ell}} \left\{ G_{0}^{\gamma \sigma}(x, y) \right\} \cdot \frac{\partial}{\partial y_{j}} \left\{ \Phi_{\varepsilon, \ell}^{*\beta \sigma}(y) \right\} \right| \\
\leq \frac{C \varepsilon \ln \left[\varepsilon^{-1} |x - y| + 2 \right]}{|x - y|^{d+1}} \tag{3.34}$$

for any $x, y \in \Omega$, where C depends only on d, m, μ , λ , τ and Ω .

Proof. Fix $x_0, y_0 \in \Omega$. Let $r = |x_0 - y_0|/8$. Since $|\nabla_x \nabla_y G_{\varepsilon}(x, y)| \leq C|x - y|^{-d}$, it suffices to consider the case $\varepsilon < r$. Fix $1 \leq \beta \leq m$ and $1 \leq j \leq d$, let

$$\begin{cases} u_{\varepsilon}^{\alpha}(x) = \frac{\partial G_{\varepsilon}^{\alpha\beta}}{\partial y_{j}}(x, y_{0}), \\ u_{0}^{\alpha}(x) = \frac{\partial}{\partial y_{j}} \left\{ \Phi_{\varepsilon, \ell}^{*\beta\sigma} \right\}(y_{0}) \cdot \frac{\partial G_{0}^{\alpha\sigma}}{\partial y_{\ell}}(x, y_{0}) \end{cases}$$

in $D_{4r} = D(x_0, 4r)$. In view of (3.24) we obtain

$$||u_{\varepsilon} - u_0||_{L^{\infty}(D_{4r})} \le C\varepsilon r^{-d} \ln \left[\varepsilon^{-1}r + 2\right].$$
 (3.35)

Since Ω is $C^{3,\eta}$, we have $\|\nabla u_0\|_{L^{\infty}(D_{4r})} \leq Cr^{-d}$,

$$\|\nabla^2 u_0\|_{L^{\infty}(D_{4r})} \le Cr^{-d-1}$$
 and $\|\nabla^2 u_0\|_{C^{0,\eta}(D_{4r})} \le Cr^{-d-1-\eta}$. (3.36)

By Lemma 3.5, estimates (3.35) and (3.36) imply that

$$\|\frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Phi_{\varepsilon,k}^{\alpha \gamma} \right\} \cdot \frac{\partial u_{0}^{\gamma}}{\partial x_{k}} \|_{L^{\infty}(D_{r})} \leq \frac{C\varepsilon \ln \left[\varepsilon^{-1}r + 2 \right]}{r^{d+1}}.$$

This gives the desired estimate (3.34).

Let $u_{\varepsilon} \in H_0^1(\Omega)$ and $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \operatorname{div}(f)$ in Ω , where $f = (f_j^{\beta}) \in L^2(\Omega)$. Then $\frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_i} = T_{\varepsilon,ij}^{\alpha\beta}(f_j^{\beta})$, where

$$T_{\varepsilon,ij}^{\alpha\beta}(g)(x) = \text{p.v.} \int_{\Omega} \frac{\partial^2}{\partial x_i \partial y_i} \left\{ G_{\varepsilon}^{\alpha\beta}(x,y) \right\} g(y) \, dy.$$

It is known that if Ω is $C^{1,\eta}$ and A(y) satisfies conditions (1.2)-(1.4), operators $T_{\varepsilon,ij}^{\alpha\beta}$ are uniformly bounded on $L^p(\Omega)$ for 1 , and of weak type (1, 1) [5].

Theorem 3.12. Suppose that Ω and A(y) satisfy the same conditions as in Theorem 3.11. Let $1 and <math>g \in L^p(\Omega)$. Then as $\varepsilon \to 0$,

$$T_{\varepsilon,ij}^{\alpha\beta}(g) - \frac{\partial}{\partial x_i} \left\{ \Phi_{\varepsilon,k}^{\alpha\gamma} \right\} \cdot T_{0,k\ell}^{\gamma\sigma} \left(\frac{\partial}{\partial x_j} \left\{ \Phi_{\varepsilon,\ell}^{*\beta\sigma} \right\} g \right) + \frac{\partial}{\partial x_i} \left\{ \Phi_{\varepsilon,k}^{\alpha\gamma} \right\} \cdot T_{0,k\ell}^{\gamma\sigma} \left(\frac{\partial}{\partial x_j} \left\{ \Phi_{\varepsilon,\ell}^{*\beta\sigma} \right\} \right) \cdot g$$

$$\rightarrow 0 \quad in \ L^q(\Omega),$$
(3.37)

for any 1 < q < p.

Proof. Let $S_{\varepsilon}(g)$ denote the left hand side of (3.37). By uniform boundedness of $\|T_{\varepsilon,ij}^{\alpha\beta}\|_{L^p\to L^p}$, $\|\nabla\Phi_{\varepsilon}\|_{\infty}$ and $\|\nabla\Phi_{\varepsilon}^*\|_{\infty}$, we see that the operators $S_{\varepsilon}:L^p(\Omega)\to L^q(\Omega)$ are uniformly bounded for $1< q< p<\infty$. As a result we may assume that $g\in C^1(\mathbb{R}^d)$.

Note that $S_{\varepsilon}(1) = 0$ and

$$S_{\varepsilon}(g)(x) = S_{\varepsilon}(g - g(x))(x) = \text{p.v.} \int_{\Omega} K_{\varepsilon}(x, y) \{g(y) - g(x)\} dy,$$

where, by Theorem 3.11 and the estimate $|\nabla_x \nabla_y G_{\varepsilon}(x,y)| \leq C|x-y|^{-d}$, the integral kernel $K_{\varepsilon}(x,y)$ satisfies

$$|K_{\varepsilon}(x,y)| \le C \min \left\{ \frac{1}{|x-y|^d}, \frac{\varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^{d+1}} \right\}.$$

It follows that if $\varepsilon < (1/2)$,

$$|S_{\varepsilon}(g)(x)| \leq C \|g\|_{C^{1}(\Omega)} \int_{B(x,\sqrt{\varepsilon})} \frac{dy}{|x-y|^{d-1}} + C\varepsilon \|g\|_{L^{\infty}(\Omega)} \int_{\Omega \setminus B(x,\sqrt{\varepsilon})} \frac{\ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^{d+1}} dy$$

$$\leq C\sqrt{\varepsilon} |\ln \varepsilon| \|g\|_{C^{1}(\Omega)}.$$

Hence, $||S_{\varepsilon}(g)||_{L^{q}(\Omega)} \to 0$ as $\varepsilon \to 0$. This completes the proof.

Finally, we consider the Dirichlet-to-Neumann map Λ_{ε} associated with the operator $\mathcal{L}_{\varepsilon}$. Let $f \in H^{1/2}(\partial\Omega)$ and $u_{\varepsilon} \in H^{1}(\Omega)$ be the solution of $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = 0$ in Ω and $u_{\varepsilon} = f$ on $\partial\Omega$. The map $\Lambda_{\varepsilon} : H^{1/2}(\partial\Omega) \to H^{-1/2}(\partial\Omega)$ is defined by $\Lambda_{\varepsilon}(f) = \frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}$. It is known that $\Lambda_{\varepsilon} : W^{1,p}(\partial\Omega) \to L^{p}(\partial\Omega)$ is uniformly bounded for $1 , if <math>\Omega$ is $C^{1,\eta}$ [21]. In the case that Ω is Lipschitz, the map is uniformly bounded for 1 if <math>m = 1 [22], and for p close to 2 if m > 1 [23]. For simplicity we will assume m = 1 and $A^* = A$ in the rest of this subsection.

Let

$$\omega_{\varepsilon}(x) = n_i(x)n_j(x)a_{ij}(x/\varepsilon) \cdot [n_k(x)n_\ell(x)\hat{a}_{k\ell}]^{-1} \cdot \frac{\partial}{\partial n} \{\Phi_{\varepsilon,s}\} \cdot n_s(x).$$

Then $\|\omega_{\varepsilon}\|_{L^{\infty}(\partial\Omega)} \leq C$. It follows from Theorem 3.11 that

$$\left| \frac{\partial}{\partial \nu_{\varepsilon}(x)} \left\{ P_{\varepsilon}(x,y) \right\} - \omega_{\varepsilon}(x) \cdot \frac{\partial}{\partial \nu_{0}(x)} \left\{ P_{0}(x,y) \right\} \cdot \omega_{\varepsilon}(y) \right| \\
\leq \frac{C \varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^{d+1}} \tag{3.38}$$

for any $x, y \in \partial \Omega$. Using $\Phi_{\varepsilon,k}(x) = x_k$ on $\partial \Omega$, one may show that $n_i n_j \hat{a}_{ij} \omega_{\varepsilon}(x) = n_k \Lambda_{\varepsilon}(x_k)$.

Theorem 3.13. Suppose that Ω and A(y) satisfy the same conditions as in Theorem 3.11. We further assume that m = 1 and $A^* = A$. Let $f \in H^1(\partial\Omega)$. Then

$$\Lambda_{\varepsilon}(f) - n_{i} \frac{\partial f}{\partial t_{ij}} \Lambda_{\varepsilon}(x_{j}) + \omega_{\varepsilon} \left[f \Lambda_{0}(\omega_{\varepsilon}) - \Lambda_{0}(\omega_{\varepsilon} f) \right]
+ \omega_{\varepsilon} n_{i} \frac{\partial f}{\partial t_{ij}} \left[\Lambda_{0}(\omega_{\varepsilon} x_{j}) - x_{j} \Lambda_{0}(\omega_{\varepsilon}) \right] \to 0,$$
(3.39)

in $L^q(\partial\Omega)$ for any 1 < q < 2, where $\frac{\partial f}{\partial t_{ij}} = (n_i \frac{\partial}{\partial x_j} - n_j \frac{\partial}{\partial x_i}) f$.

Proof. By a linear change of variables Theorems 5.1 and 5.2 in Section 5 continue to hold for $\Lambda = \Lambda_0$. It follows that $\|\Lambda_0(\omega_{\varepsilon}f) - f\Lambda_0(\omega_{\varepsilon})\|_{L^2(\partial\Omega)} \leq C\|f\|_{H^1(\partial\Omega)}$ and $\|\Lambda_0(\omega_{\varepsilon}x_j) - x_j\Lambda_0(\omega_{\varepsilon})\|_{L^p(\partial\Omega)} \leq C_p\|\omega_{\varepsilon}\|_{L^p(\partial\Omega)} \leq C_p$ for any $1 . Hence, by Hölder's inequality, the left hand side of (3.39), as an operator, is uniformly bounded from <math>H^1(\partial\Omega)$ to $L^q(\partial\Omega)$ for any 1 < q < 2. Consequently, it suffices to show that for $f \in C^2(\partial\Omega)$, the left hand side of (3.39) goes to zero in $L^{\infty}(\partial\Omega)$, as $\varepsilon \to 0$.

To this end, recall that $\Lambda_{\varepsilon}(f) = \frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}$, where $u_{\varepsilon}(z) = \int_{\Omega} P_{\varepsilon}(z, y) f(y) d\sigma(y)$ for $z \in \Omega$. Write

$$u_{\varepsilon}(z) - f(x) = \int_{\partial\Omega} P_{\varepsilon}(z, y) \left\{ f(y) - f(x) \right\} d\sigma(y)$$

$$= \int_{\partial\Omega} P_{\varepsilon}(z, y) \left\{ f(y) - f(x) - n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \cdot (y_{j} - x_{j}) \right\} d\sigma(y)$$

$$+ \int_{\partial\Omega} P_{\varepsilon}(z, y) n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \cdot (y_{j} - x_{j}) d\sigma(y)$$

$$= \int_{\partial\Omega} P_{\varepsilon}(z, y) \left\{ f(y) - f(x) - n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \cdot (y_{j} - x_{j}) \right\} d\sigma(y)$$

$$+ n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \cdot \left[\Phi_{\varepsilon, j}(z) - x_{j} \right].$$
(3.40)

Since $|f(y) - f(x) - n_i(x) \cdot \partial f / \partial t_{ij}(x) \cdot (y_j - x_j)| \le C_f |y - x|^2$ for $x, y \in \partial \Omega$, it follows by taking derivatives in z and then letting $z \to x$ in (3.40) that

$$\Lambda_{\varepsilon}(f)(x) = \int_{\partial\Omega} \frac{\partial}{\partial\nu_{\varepsilon}(x)} \left\{ P_{\varepsilon}(x,y) \right\} \left\{ f(y) - f(x) - n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \cdot (y_{j} - x_{j}) \right\} d\sigma(y) + n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \Lambda_{\varepsilon}(x_{j}).$$

In view of (3.38) as well as the estimate $|\nabla_x P_{\varepsilon}(x,y)| \leq C|\nabla_x \nabla_y G_{\varepsilon}(x,y)| \leq C|x-y|^{-d}$, we obtain

$$\Lambda_{\varepsilon}(f)(x) = I_{\varepsilon}(x) + n_{i} \frac{\partial f}{\partial t_{ij}} \Lambda_{\varepsilon}(x_{j})
+ \int_{\partial \Omega} \omega_{\varepsilon}(x) \frac{\partial}{\partial \nu_{0}(x)} \left\{ P_{0}(x, y) \right\} \omega_{\varepsilon}(y) \left\{ f(y) - f(x) - n_{i}(x) \frac{\partial f}{\partial t_{ij}(x)} \cdot (y_{j} - x_{j}) \right\} d\sigma(y)
= I_{\varepsilon}(x) + n_{i} \frac{\partial f}{\partial t_{ij}} \Lambda_{\varepsilon}(x_{j}) + \omega_{\varepsilon} \left[\Lambda_{0}(\omega_{0}f) - f \Lambda_{0}(\omega_{\varepsilon}) \right] + \omega_{\varepsilon} n_{i} \frac{\partial f}{\partial t_{ij}} \left[x_{j} \Lambda_{0}(\omega_{\varepsilon}) - \Lambda_{0}(\omega_{\varepsilon}x_{j}) \right],$$

where the term $I_{\varepsilon}(x)$ satisfies

$$|I_{\varepsilon}(x)| \leq C_f \left\{ \varepsilon \int_{\partial \Omega \setminus B(x,\varepsilon)} \frac{\ln[\varepsilon^{-1}|x-y|+2]}{|y-x|^{d-1}} d\sigma(y) + \int_{B(x,\varepsilon)\cap\partial\Omega} \frac{d\sigma(y)}{|y-x|^{d-2}} \right\}$$

$$\leq C_f \varepsilon [\ln(\varepsilon^{-1}M+2)]^2.$$

This gives the desired estimate.

4 Asymptotic behavior of Neumann functions

Throughout this section we will assume that A(y) satisfies conditions (1.2), (1.3) and (1.4). Under these conditions one may construct a matrix of Neumann functions $N_{\varepsilon}(x,y) = (N_{\varepsilon}^{\alpha\beta}(x,y))$ in a bounded Lipschitz domain Ω such that

$$\begin{cases}
\mathcal{L}_{\varepsilon} \{ N_{\varepsilon}^{\beta}(\cdot, y) \} = e^{\beta} \delta_{y}(x) & \text{in } \Omega, \\
\frac{\partial}{\partial \nu_{\varepsilon}} \{ N_{\varepsilon}^{\beta}(\cdot, y) \} = -e^{\beta} |\partial \Omega|^{-1} & \text{on } \partial \Omega, \\
\int_{\partial \Omega} N_{\varepsilon}^{\beta}(x, y) d\sigma(x) = 0,
\end{cases} (4.1)$$

where $N_{\varepsilon}^{\beta}(x,y) = (N_{\varepsilon}^{1\beta}(x,y), \dots, N_{\varepsilon}^{m\beta}(x,y))$ and $e^{\beta} = (0,\dots,1,\dots,0)$ with 1 in the β^{th} position. Let $u_{\varepsilon} \in H^{1}(\Omega)$ be a solution to $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω and $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = g$ on $\partial \Omega$. Then

$$u_{\varepsilon}(x) - \frac{1}{|\partial\Omega|} \int_{\partial\Omega} u_{\varepsilon} = \int_{\Omega} N_{\varepsilon}(x, y) F(y) \, dy + \int_{\partial\Omega} N_{\varepsilon}(x, y) g(y) \, d\sigma(y). \tag{4.2}$$

Under the additional assumption that $A^* = A$ and Ω is $C^{1,\eta}$ for some $\eta \in (0,1)$, it was proved in [21] that

$$\begin{cases}
|N_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^{d-2}}, \\
|\nabla_x N(x,y)| + |\nabla_y N_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^{d-1}}, \\
|\nabla_x \nabla_y N_{\varepsilon}(x,y)| \leq \frac{C}{|x-y|^d}
\end{cases} (4.3)$$

for any $x, y \in \Omega$. The goal of this section is to establish the asymptotic estimates of $N_{\varepsilon}(x, y)$ and $\nabla_x N_{\varepsilon}(x, y)$ in Theorem 1.2.

4.1 L^{∞} estimates

The goal of this subsection is to prove the estimate (1.9). We also obtain an $O(\varepsilon)$ estimate (up to a logarithmic factor) for $||u_{\varepsilon}-u_0||_{L^p(\Omega)}$ for solutions with Neumann conditions. Recall that $D_r = D_r(x_0, r) = B(x_0, r) \cap \Omega$ and $\Delta_r = \Delta(x_0, r) = B(x_0, r) \cap \partial\Omega$ for some $x_0 \in \overline{\Omega}$ and $0 < r < r_0$. We begin with an L^{∞} estimate for local solutions.

Lemma 4.1. Let Ω be a bounded $C^{1,\eta}$ domain for some $\eta \in (0,1)$. Let $u_{\varepsilon} \in H^1(D_{3r})$ and $u_0 \in W^{2,p}(D_{3r})$ for some p > d. Suppose that $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in D_{3r} and $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\partial u_0}{\partial \nu_0}$ on Δ_{3r} . Then, if $0 < \varepsilon < (r/2)$,

$$||u_{\varepsilon} - u_{0}||_{L^{\infty}(D_{r})} \leq \frac{C}{r^{d}} \int_{D_{3r}} |u_{\varepsilon} - u_{0}| + C\varepsilon \ln[\varepsilon^{-1}r + 2] ||\nabla u_{0}||_{L^{\infty}(D_{3r})} + C_{p} \varepsilon r^{1 - \frac{d}{p}} ||\nabla^{2}u_{0}||_{L^{p}(D_{3r})}.$$

$$(4.4)$$

Proof. By rescaling we may assume r=1. Choose a $C^{1,\eta}$ domain \widetilde{D} such that $D_2\subset \widetilde{D}\subset D_3$. Let

$$w_{\varepsilon} = u_{\varepsilon}(x) - u_0(x) - \varepsilon \chi_j^{\beta}(x/\varepsilon) \frac{\partial u_0^{\beta}}{\partial x_i}.$$

Using Proposition 2.2 with $V_{\varepsilon,j}^{\beta} = P_j^{\beta}(x) + \varepsilon \chi_j^{\beta}(x/\varepsilon)$, we see that

$$(\mathcal{L}(w_{\varepsilon}))^{\alpha} = \varepsilon \frac{\partial}{\partial x_{i}} \left\{ b_{ijk}^{\alpha\beta}(x/\varepsilon) \frac{\partial^{2} u_{0}^{\beta}}{\partial x_{j} \partial x_{k}} \right\}, \tag{4.5}$$

where $b_{ijk}^{\alpha\beta}(y) = F_{jik}^{\alpha\beta}(y) + a_{ij}^{\alpha\gamma}(y)\chi_k^{\gamma\beta}(y)$ is a bounded periodic function. Also, by a direct computation, we have

$$\left(\frac{\partial w_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} = \left(\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} - \left(\frac{\partial u_{0}}{\partial \nu_{0}}\right)^{\alpha} + \frac{\varepsilon}{2} \left(n_{i} \frac{\partial}{\partial x_{j}} - n_{j} \frac{\partial}{\partial x_{i}}\right) \left\{F_{jik}^{\alpha\gamma}(x/\varepsilon) \frac{\partial u_{0}^{\gamma}}{\partial x_{k}}\right\} - \varepsilon n_{i} b_{ijk}^{\alpha\beta}(x/\varepsilon) \frac{\partial^{2} u_{0}^{\beta}}{\partial x_{j} \partial x_{k}} \tag{4.6}$$

(see [20, Lemma 5.1]). Let $w_{\varepsilon} = w_{\varepsilon}^{(1)} + w_{\varepsilon}^{(2)}$, where

$$\left(w_{\varepsilon}^{(1)}(x)\right)^{\alpha} = -\varepsilon \int_{\widetilde{D}} \frac{\partial}{\partial y_i} \left\{ \widetilde{N}_{\varepsilon}^{\alpha\beta}(x,y) \right\} \cdot b_{ijk}^{\beta\gamma}(y/\varepsilon) \cdot \frac{\partial^2 u_0^{\gamma}}{\partial y_j \partial y_k} \, dy \tag{4.7}$$

and $\widetilde{N}_{\varepsilon}(x,y)$ denotes the matrix of Neumann functions for $\mathcal{L}_{\varepsilon}$ in \widetilde{D} . Since $|\nabla_{y}\widetilde{N}_{\varepsilon}(x,y)| \leq C|x-y|^{1-d}$, it follows from Hölder's inequality that

$$||w_{\varepsilon}^{(1)}||_{L^{\infty}(D_2)} \le C_p \varepsilon ||\nabla^2 u_0||_{L^p(D_3)} \qquad \text{for any } p > d.$$

$$\tag{4.8}$$

To estimate $w_{\varepsilon}^{(2)}$, we observe that $\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(2)}) = 0$ in \widetilde{D} and

$$\left(\frac{\partial w_{\varepsilon}^{(2)}}{\partial \nu_{\varepsilon}}\right)^{\alpha} = \left(\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} - \left(\frac{\partial u_{0}}{\partial \nu_{0}}\right)^{\alpha} + \frac{\varepsilon}{2} \left(n_{i} \frac{\partial}{\partial x_{j}} - n_{j} \frac{\partial}{\partial x_{i}}\right) \left\{F_{jik}^{\alpha \gamma}(x/\varepsilon) \frac{\partial u_{0}^{\gamma}}{\partial x_{k}}\right\} \tag{4.9}$$

on $\partial \widetilde{D}$. Let $w_{\varepsilon}^{(2)} = w_{\varepsilon}^{(21)} + w_{\varepsilon}^{(22)}$, where

$$(w_{\varepsilon}^{(21)})^{\alpha}(x) = -\frac{\varepsilon}{2} \int_{\partial \widetilde{D}} \left(n_i \frac{\partial}{\partial y_i} - n_j \frac{\partial}{\partial y_i} \right) \left\{ \widetilde{N}_{\varepsilon}^{\alpha\beta}(x, y) \right\} \cdot F_{jik}^{\beta\gamma}(y/\varepsilon) \frac{\partial u_0^{\gamma}}{\partial y_k} d\sigma(y).$$
 (4.10)

For each $x \in \widetilde{D}$, choose $\hat{x} \in \partial \widetilde{D}$ such that $|\hat{x} - x| = \operatorname{dist}(x, \partial \widetilde{D})$. Note that for $y \in \partial \widetilde{D}$, $|y - \hat{x}| \leq |y - x| + |x - \hat{x}| \leq 2|y - x|$. Thus, $|\nabla_y \widetilde{N}_{\varepsilon}(x, y)| \leq C|y - \hat{x}|^{1-d}$ and

$$\begin{split} |\left(w_{\varepsilon}^{(21)}\right)^{\alpha}(x)| &= \frac{\varepsilon}{2} \Big| \int_{\partial \widetilde{D}} \left(n_{i} \frac{\partial}{\partial y_{j}} - n_{j} \frac{\partial}{\partial y_{i}} \right) \left\{ \widetilde{N}_{\varepsilon}^{\alpha\beta}(x, y) \right\} \cdot \left\{ f_{ji}^{\beta}(y) - f_{ji}^{\beta}(\hat{x}) \right\} d\sigma(y) \Big| \\ &\leq C \varepsilon \int_{\partial \widetilde{D}} \frac{|f(y) - f(\hat{x})|}{|y - \hat{x}|^{d-1}} d\sigma(y), \end{split}$$

where $f(y) = (f_{ji}^{\beta}(y)) = (F_{jik}^{\beta\gamma}(y/\varepsilon)\frac{\partial u_0^{\gamma}}{\partial y_k}(y))$. Since $||f||_{L^{\infty}(D_3)} \le C||\nabla u_0||_{L^{\infty}(D_3)}$ and

$$|f(y) - f(\hat{x})| \le C\varepsilon^{-1}|y - \hat{x}|\|\nabla u_0\|_{L^{\infty}(D_3)} + |y - \hat{x}|^{\rho}\|\nabla u_0\|_{C^{0,\rho}(D_3)},$$

where $0 < \rho < \eta$ and we have used the fact $||F_{jik}^{\beta\gamma}||_{C^1(Y)} \leq C$, it follows that

$$|w_{\varepsilon}^{(21)}(x)| \leq C\varepsilon \|\nabla u_{0}\|_{L^{\infty}(D_{3})} \int_{\partial \widetilde{D} \setminus B(\hat{x},\varepsilon)} |\hat{x} - y|^{1-d} d\sigma(y)$$

$$+ C \|\nabla u_{0}\|_{L^{\infty}(D_{3})} \int_{B(\hat{x},\varepsilon)\cap\partial \widetilde{D}} |y - \hat{x}|^{2-d} d\sigma(y)$$

$$+ C\varepsilon \|\nabla u_{0}\|_{C^{0,\rho}(D_{3})} \int_{B(\hat{x},\varepsilon)\cap\partial \widetilde{D}} |y - \hat{x}|^{1-d+\rho} d\sigma(y)$$

$$\leq C\varepsilon \ln[\varepsilon^{-1} + 2] \|\nabla u_{0}\|_{L^{\infty}(D_{3})} + C\varepsilon^{1+\rho} \|\nabla u_{0}\|_{C^{0,\rho}(D_{3})}.$$

By Sobolev imbedding, this implies that

$$||w_{\varepsilon}^{(21)}||_{L^{\infty}(D_{2})} \le C\varepsilon \ln[\varepsilon^{-1} + 2]||\nabla u_{0}||_{L^{\infty}(D_{3})} + C_{p}\varepsilon||\nabla^{2}u_{0}||_{L^{p}(D_{3})}$$
(4.11)

for any p > d.

Finally, to estimate $w_{\varepsilon}^{(22)}$, we note that $\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(22)}) = 0$ in D_2 and $\frac{\partial}{\partial \nu_{\varepsilon}} \{w_{\varepsilon}^{(22)}\} = \frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} - \frac{\partial u_0}{\partial \nu_0} = 0$ on Δ_2 . It follows from [21, Theorem 3.1] that

$$||w_{\varepsilon}^{(22)}||_{L^{\infty}(D_{1})} \leq C \int_{D_{2}} |w_{\varepsilon}^{(22)}| dx$$

$$\leq C \int_{D_{2}} |u_{\varepsilon} - u_{0}| dx + C\varepsilon ||\nabla u_{0}||_{L^{\infty}(D_{2})}$$

$$+ C||w_{\varepsilon}^{(1)}||_{L^{\infty}(D_{2})} + C||w_{\varepsilon}^{(21)}||_{L^{\infty}(D_{2})}.$$

Hence,

$$||u_{\varepsilon} - u_{0}||_{L^{\infty}(D_{1})} \leq C \int_{D_{2}} |u_{\varepsilon} - u_{0}| dx + C\varepsilon ||\nabla u_{0}||_{L^{\infty}(D_{2})} + C||w_{\varepsilon}^{(1)}||_{L^{\infty}(D_{2})} + C||w_{\varepsilon}^{(21)}||_{L^{\infty}(D_{2})}.$$

This, together with (4.8) and (4.11), gives (4.4).

The next lemma on the traces of fractional integrals is known. We provide a proof for the sake of completness.

Lemma 4.2. Let Ω be a bounded Lipschitz domain. Let $g \in C_0^1(\Omega \cap B(y_0, r))$ for some $y_0 \in \overline{\Omega}$ and

$$u(x) = \int_{\Omega} \frac{g(y) \, dy}{|x - y|^{d-1}}.$$

Then $||u||_{L^2(\partial\Omega)} \le Cr^{1/2}||g||_{L^2(\Omega)}$.

Proof. By the well known estimates for fractional and singular integrals, $\|\nabla u\|_{L^2(\Omega)} \le C\|g\|_{L^2(\Omega)}$ and $\|u\|_{L^2(\Omega)} \le C\|g\|_{L^p(\Omega)}$, where $p = \frac{2d}{d+2}$. Choose a vector field $v \in C_0^{\infty}(\mathbb{R}^d, \mathbb{R}^d)$ such that $\langle v, n \rangle \ge c_0 > 0$ on $\partial\Omega$ and $|v| \le 1$, $|\nabla v| \le C/r$ in \mathbb{R}^d . It follows from integration by parts and Hölder's inequality that

$$c_{0} \int_{\partial\Omega} |u|^{2} d\sigma \leq \int_{\partial\Omega} \langle v, n \rangle |u|^{2} d\sigma \leq Cr^{-1} \int_{\Omega} |u|^{2} dx + C \int_{\Omega} |u| |\nabla u| dx$$

$$\leq Cr^{-1} \int_{\Omega} |u|^{2} dx + Cr \int_{\Omega} |\nabla u|^{2} dx$$

$$\leq Cr^{-1} ||g||_{L^{2}(\Omega)}^{2} + Cr ||g||_{L^{2}(\Omega)}^{2}$$

$$\leq Cr ||g||_{L^{2}(\Omega)}^{2},$$

where we have used the fact $supp(g) \subset B(y_0, r)$. This completes the proof.

We are now ready to prove the estimate (1.9).

Theorem 4.3. Suppose that A(y) satisfies conditions (1.2), (1.3), (1.4) and $A^* = A$. Let Ω be a bounded $C^{1,1}$ domain. Then

$$|N_{\varepsilon}(x,y) - N_0(x,y)| \le \frac{C\varepsilon \ln[\varepsilon^{-1}|x-y|+2]}{|x-y|^{d-1}},\tag{4.12}$$

for any $x, y \in \Omega$, where C depends only on d, m, μ , λ , τ and Ω .

Proof. By rescaling we may assume that $\operatorname{diam}(\Omega) = 1$. Fix $x_0, y_0 \in \Omega$ and let $r = |x_0 - y_0|/8$. Since $|N_{\varepsilon}(x_0, y_0)| \leq Cr^{2-d}$, we may assume that $\varepsilon < r$. For $g \in C_0^{\infty}(D(y_0, r))$ and $\varepsilon \geq 0$, let $u_{\varepsilon}(x) = \int_{\Omega} N_{\varepsilon}(x, y) g(y) \, dy$. Then $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = g$ in Ω , $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = -\frac{1}{|\partial \Omega|} \int_{\Omega} g$ on $\partial \Omega$ and $\int_{\partial \Omega} u_{\varepsilon} = 0$. It follows that $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in Ω and $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\partial u_0}{\partial \nu_0}$ on $\partial \Omega$. Let $w_{\varepsilon} = u_{\varepsilon}(x) - u_0(x) - \varepsilon \chi_j^{\beta}(x/\varepsilon) \frac{\partial u_0^{\beta}}{\partial x_j}$. As in the proof of Lemma 4.1,

$$(\mathcal{L}(w_{\varepsilon}))^{\alpha} = \varepsilon \frac{\partial}{\partial x_i} \left\{ b_{ijk}^{\alpha\beta}(x/\varepsilon) \frac{\partial^2 u_0^{\beta}}{\partial x_j \partial x_k} \right\} \quad \text{in } \Omega$$

and

$$\left(\frac{\partial w_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} = \frac{\varepsilon}{2} \left(n_{i} \frac{\partial}{\partial x_{j}} - n_{j} \frac{\partial}{\partial x_{i}}\right) \left\{F_{jik}^{\alpha \gamma}(x/\varepsilon) \frac{\partial u_{0}^{\gamma}}{\partial x_{k}}\right\} - \varepsilon n_{i} b_{ijk}^{\alpha \beta}(x/\varepsilon) \frac{\partial^{2} u_{0}^{\beta}}{\partial x_{j} \partial x_{k}}$$

on $\partial\Omega$. Now, write $w_{\varepsilon} = \theta_{\varepsilon} + z_{\varepsilon} + \rho$, where $\mathcal{L}_{\varepsilon}(z_{\varepsilon}) = \mathcal{L}_{\varepsilon}(w_{\varepsilon})$ in Ω , $\int_{\Omega} z_{\varepsilon} = 0$,

$$\frac{\partial z_{\varepsilon}}{\partial \nu_{\varepsilon}} = -\varepsilon n_i b_{ijk}^{\alpha\beta}(x/\varepsilon) \frac{\partial^2 u_0^{\beta}}{\partial x_j \partial x_k} \quad \text{on } \partial \Omega,$$

and $\rho = \frac{1}{|\partial\Omega|} \int_{\partial\Omega} \left\{ w_{\varepsilon} - z_{\varepsilon} \right\}$ is a constant. Note that $\|\nabla z_{\varepsilon}\|_{L^{2}(\Omega)} \leq C\varepsilon \|\nabla^{2}u_{0}\|_{L^{2}(\Omega)} \leq C\varepsilon \|g\|_{L^{2}(\Omega)}$. Since $\int_{\Omega} z_{\varepsilon} = 0$, by the Poincaré inequality, we obtain $\|z_{\varepsilon}\|_{L^{p}(\Omega)} \leq C\varepsilon \|g\|_{L^{2}(\Omega)}$, where $p = \frac{2d}{d-2}$. It follows by Hölder's inequality that

$$||z_{\varepsilon}||_{L^{2}(D(x_{0},r))} \leq Cr^{\frac{d}{2}-\frac{d}{p}}||z_{\varepsilon}||_{L^{p}(D(x_{0},r))} \leq C\varepsilon r||g||_{L^{2}(\Omega)}.$$
(4.13)

Next, to estimate θ_{ε} , we observe that $\mathcal{L}_{\varepsilon}(\theta_{\varepsilon}) = 0$ in Ω , $\int_{\partial\Omega} \theta_{\varepsilon} = 0$ and

$$\frac{\partial \theta_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\varepsilon}{2} \left(n_i \frac{\partial}{\partial x_j} - n_j \frac{\partial}{\partial x_i} \right) \left\{ F_{jik}^{\alpha \gamma}(x/\varepsilon) \frac{\partial u_0^{\gamma}}{\partial x_k} \right\}.$$

Using a duality argument and L^2 estimates for the Neumann problem in [22], we may deduce that $\|\theta_{\varepsilon}\|_{L^2(\partial\Omega)} \leq C\varepsilon \|\nabla u_0\|_{L^2(\partial\Omega)}$ (see the proof of Theorem 5.2 in [20]). By the square function estimate for the L^2 Dirichlet problem [20], this implies that

$$\|\theta_{\varepsilon}\|_{H^{1/2}(\Omega)} \le C\|\theta_{\varepsilon}\|_{L^{2}(\partial\Omega)} \le C\varepsilon \|\nabla u_{0}\|_{L^{2}(\partial\Omega)}.$$

It follows from the Sobolev imbedding that $\|\theta_{\varepsilon}\|_{L^{p_1}(\Omega)} \leq C\varepsilon \|\nabla u_0\|_{L^2(\partial\Omega)}$, where $p_1 = \frac{2d}{d-1}$. By Hölder's inequality, this gives

$$\|\theta_{\varepsilon}\|_{L^{2}(D(x_{0},r))} \leq C\varepsilon r^{\frac{1}{2}} \|\nabla u_{0}\|_{L^{2}(\partial\Omega)}. \tag{4.14}$$

Since

$$|\nabla u_0(x)| \le C \int_{\Omega} \frac{|g(y)| \, dy}{|x - y|^{d-1}},$$

we may invoke Lemma 4.2 to claim that $\|\nabla u_0\|_{L^2(\Omega)} \leq Cr^{1/2} \|g\|_{L^2(\Omega)}$. In view of (4.14) we obtain $\|\theta_{\varepsilon}\|_{L^2(D(x_0,r))} \leq C\varepsilon r \|g\|_{L^2(\Omega)}$. This, together with (4.13) and the observation

$$|\rho| \le C \int_{\partial\Omega} \{ \varepsilon |\nabla u_0| + |z_{\varepsilon}| \} d\sigma \le C \varepsilon ||g||_{L^2(\Omega)},$$

gives $||w_{\varepsilon}||_{L^{2}(D(x_{0},r))} \leq C\varepsilon r||g||_{L^{2}(\Omega)}$. It follows that

$$\left\{ \frac{1}{r^d} \int_{D(x_0,r)} |u_{\varepsilon} - u_0|^2 \right\}^{1/2} \le C \varepsilon r^{\frac{2-d}{2}} ||g||_{L^2(\Omega)}.$$

Since $\|\nabla u_0\|_{L^{\infty}(D(x_0,r))} \leq Cr^{\frac{2-d}{2}}\|g\|_{L^2(\Omega)}$ and $\|\nabla^2 u_0\|_{L^p(\Omega)} \leq C\|g\|_{L^p(\Omega)}$, by Lemma 4.1, we obtain

$$|u_{\varepsilon}(x_0) - u_0(x_0)| \le C_p \varepsilon r^{1-\frac{d}{p}} \ln[\varepsilon^{-1}r + 2] ||g||_{L^p(\Omega)},$$

where p > d. By duality this gives

$$\left\{ \frac{1}{r^d} \int_{D(y_0,r)} |N_{\varepsilon}(x_0,y) - N_0(x_0,y)|^{p'} dy \right\}^{1/p'} \le C_p \, \varepsilon r^{1-d} \ln[\varepsilon^{-1}r + 2].$$

Finally, since

$$\frac{\partial}{\partial \nu_{\varepsilon}(y)} \{ N_{\varepsilon}(x,y) \} = \frac{\partial}{\partial \nu_{0}(y)} \{ N_{0}(x,y) \} = -\frac{1}{|\partial \Omega|} \quad \text{on } \partial \Omega$$

 $|\nabla_y N_0(x,y)| \le C|x-y|^{1-d}$ and $\|\nabla_y^2 N_0(x_0,y)\|_{L^p(D(y_0,r))} \le Cr^{\frac{d}{p}-d}$, we may invoke Lemma 4.1 again to obtain

$$|N_{\varepsilon}(x_0, y_0) - N_0(x_0, y_0)| \leq \frac{C}{r^d} \int_{D(y_0, r)} |N_{\varepsilon}(x_0, y) - N_0(x_0, y)| dy$$
$$+ C\varepsilon r^{1-d} \ln[\varepsilon^{-1}r + 2]$$
$$\leq C\varepsilon r^{1-d} \ln[\varepsilon^{-1}r + 2].$$

This completes the proof.

Remark 4.4. It is not clear whether the logarithmic factor in (4.12) is necessary. Also, in view of Theorem 3.3 on Green's functions, it would be interesting to show that the estimate (4.12) holds for scalar equations with no smoothness condition on the coefficients.

As a corollary of Theorem 4.3, we obtain an $O(\varepsilon)$ (up to a logarithmic factor) estimate for $||u_{\varepsilon} - u_0||_{L^q(\Omega)}$.

Theorem 4.5. Suppose that A(y) and Ω satisfy the same assumptions as in Theorem 4.3. Let $1 . For <math>\varepsilon \ge 0$ and $F \in L^p(\Omega)$ with $\int_{\Omega} F = 0$, let $u_{\varepsilon} \in W^{1,p}(\Omega)$ be the solution to the Neumann problem: $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω , $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = 0$ and $\int_{\partial \Omega} u_{\varepsilon} = 0$. Then

$$||u_{\varepsilon} - u_0||_{L^q(\Omega)} \le C\varepsilon \ln[\varepsilon^{-1}M + 2]||F||_{L^p(\Omega)}$$
(4.15)

holds if $1 and <math>\frac{1}{q} = \frac{1}{p} - \frac{1}{d}$, or p > d and $q = \infty$, where $M = \operatorname{diam}(\Omega)$. Moreover,

$$||u_{\varepsilon} - u_0||_{L^{\infty}(\Omega)} \le C\varepsilon \left[\ln(\varepsilon^{-1}M + 2)\right]^{2 - \frac{1}{d}} ||F||_{L^d(\Omega)}. \tag{4.16}$$

Proof. Note that by the estimate (4.12),

$$|u_{\varepsilon}(x) - u_{0}(x)| \leq \int_{\Omega} |N_{\varepsilon}(x, y) - N_{0}(x, y)| |F(y)| \, dy$$
$$\leq C\varepsilon \ln(\varepsilon^{-1}M + 2) \int_{\Omega} \frac{|F(y)| \, dy}{|x - y|^{d-1}}.$$

The rest of the proof is the same as that of Theorem 3.4.

4.2 Lipschitz estimates

In this subsection we give the proof of the estimate (1.10). We also establish an $O(\varepsilon^t)$ estimate for $u_{\varepsilon} - u_0 - \{\Psi_{\varepsilon,j} - P_j\}\frac{\partial u_0}{\partial x_j}$ in $W^{1,p}(\Omega)$ for any $t \in (0,1)$ and 1 .

Lemma 4.6. Let Ω be a bounded $C^{1,\eta}$ domain for some $\eta \in (0,1]$ and

$$u_{\varepsilon}(x) = \text{p.v.} \int_{\Omega} \frac{\partial^2}{\partial x_i \partial y_j} \{ N_{\varepsilon}(x, y) \} f(y) \, dy$$

for some $1 \le i, j \le d$. Then

$$||u_{\varepsilon}||_{L^{\infty}(\Omega)} \le C \left\{ \ln[\varepsilon^{-1}M + 2] + M^{\eta} \right\} ||f||_{L^{\infty}(\Omega)} + C\varepsilon^{\eta} H_{\varepsilon,\eta}(f), \tag{4.17}$$

where $M = \operatorname{diam}(\Omega)$ and $H_{\varepsilon,\eta}(f) = \sup \left\{ \frac{|f(x) - f(y)|}{|x - y|^{\eta}} : x, y \in \Omega \text{ and } |x - y| < \varepsilon \right\}.$

Proof. For $x \in \Omega$, choose $\hat{x} \in \partial \Omega$ such that $|x - \hat{x}| = \operatorname{dist}(x, \partial \Omega)$. Note that

$$u_{\varepsilon}(x) = \int_{\Omega} \frac{\partial^{2}}{\partial x_{i} \partial y_{j}} \left\{ N_{\varepsilon}(x, y) \right\} \cdot \left\{ f(y) - f(x) \right\} dy$$
$$+ \int_{\partial \Omega} \left\{ n_{j}(y) - n_{j}(\hat{x}) \right\} \cdot \frac{\partial}{\partial x_{i}} \left\{ N_{\varepsilon}(x, y) \right\} \cdot f(x) d\sigma(y),$$

where we have used the fact $\int_{\partial\Omega} N_{\varepsilon}(x,y) d\sigma(y) = 0$. This, together with the estimates $|\nabla_x N_{\varepsilon}(x,y)| \leq C|x-y|^{1-d}$ and $|\nabla_x \nabla_y N_{\varepsilon}(x,y)| \leq C|x-y|^{-d}$, gives

$$||u_{\varepsilon}||_{L^{\infty}(\Omega)} \leq C \int_{\Omega} \frac{|f(y) - f(x)|}{|x - y|^d} dy + C||f||_{L^{\infty}(\Omega)} \int_{\partial \Omega} \frac{d\sigma(y)}{|y - \hat{x}|^{d-1-\eta}}$$

$$\leq C||f||_{L^{\infty}(\Omega)} \int_{\Omega \setminus B(x,\varepsilon)} \frac{dy}{|x - y|^d} + CH_{\varepsilon,\eta}(f) \int_{|y - x| < \varepsilon} \frac{dy}{|y - x|^{d-\eta}} + C||f||_{L^{\infty}(\Omega)} M^{\eta}$$

$$\leq C||f||_{L^{\infty}(\Omega)} \ln[\varepsilon^{-1}M + 2] + C\varepsilon^{\eta} H_{\varepsilon,\eta}(f) + C||f||_{L^{\infty}(\Omega)} M^{\eta}.$$

This completes the proof.

The following lemma provides a Lipschitz estimate for local solutions with Neumann boundary conditions.

Lemma 4.7. Let Ω be a bounded $C^{2,\eta}$ domain for some $\eta \in (0,1)$. Suppose that $u_{\varepsilon} \in H^1(D_{3r})$, $u_0 \in C^{2,\eta}(D_{3r})$ and $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_0(u_0)$ in D_{3r} . Also assume that $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\partial u_0}{\partial \nu_0}$ on Δ_{3r} . Then, if $0 < \varepsilon < (r/2)$,

$$\|\frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Psi_{\varepsilon,j}^{\alpha\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \|_{L^{\infty}(D_{r})}$$

$$\leq \frac{C}{r^{d+1}} \int_{D_{3r}} |u_{\varepsilon} - u_{0}| \, dx + C\varepsilon r^{-1} \ln[\varepsilon^{-1}M + 2] \|\nabla u_{0}\|_{L^{\infty}(D_{3r})}$$

$$+ C\varepsilon^{1-\rho} r^{\rho} \ln[\varepsilon^{-1}M + 2] \|\nabla^{2}u_{0}\|_{L^{\infty}(D_{3r})} + C\varepsilon r^{\rho} \ln[\varepsilon^{-1}M + 2] \|\nabla^{2}u_{0}\|_{C^{0,\rho}(D_{3r})}$$

$$(4.18)$$

for any $0 < \rho < \min(\eta, \tau)$, where $(\Psi_{\varepsilon,j}^{\beta})$ denotes the matrix of Neumann correctors for $\mathcal{L}_{\varepsilon}$ in Ω and $M = \operatorname{diam}(\Omega)$.

Proof. By rescaling and translation we may assume that r=1 and $0 \in D_1$. Let

$$w_{\varepsilon} = u_{\varepsilon}(x) - u_0(x) - \left\{ \Psi_{\varepsilon,j}^{\beta} - P_j^{\beta} \right\} \cdot \frac{\partial u_0^{\beta}}{\partial x_i}.$$

Choose a $C^{2,\eta}$ domain such that $D_2 \subset \widetilde{D} \subset D_3$. We now write

$$w_{\varepsilon}(x) = \int_{\widetilde{D}} \widetilde{N}_{\varepsilon}(x, y) \mathcal{L}_{\varepsilon}(w_{\varepsilon}) dy + \int_{\partial \widetilde{D}} \widetilde{N}_{\varepsilon}(x, y) \frac{\partial w_{\varepsilon}}{\partial \nu_{\varepsilon}} d\sigma(y) + \frac{1}{|\partial \widetilde{D}|} \int_{\partial \widetilde{D}} w_{\varepsilon}$$

for $x \in D_2$, where $\widetilde{N}_{\varepsilon}(x,y)$ denotes the matrix of Neumann functions for $\mathcal{L}_{\varepsilon}$ in \widetilde{D} . In view of Propositions 2.2 and 2.3, we have $w_{\varepsilon} = w_{\varepsilon}^{(1)} + w_{\varepsilon}^{(2)} + c$, where $c = \frac{1}{|\partial \widetilde{D}|} \int_{\partial \widetilde{D}} w_{\varepsilon}$,

$$w_{\varepsilon}^{(1)}(x) = -\varepsilon \int_{\widetilde{D}} \frac{\partial}{\partial y_{i}} \left\{ \widetilde{N}_{\varepsilon}(x, y) \right\} \cdot \left\{ F_{jik}(y/\varepsilon) \right\} \cdot \frac{\partial^{2} u_{0}}{\partial y_{j} \partial y_{k}} dy$$

$$- \int_{\widetilde{D}} \frac{\partial}{\partial y_{i}} \left\{ \widetilde{N}_{\varepsilon}(x, y) \right\} \cdot a_{ij}(y/\varepsilon) \left\{ \Psi_{\varepsilon, k}(y) - P_{k}(y) \right\} \cdot \frac{\partial^{2} u_{0}}{\partial y_{j} \partial y_{k}} dy$$

$$+ \int_{\widetilde{D}} \widetilde{N}_{\varepsilon}(x, y) \cdot a_{ij}(y/\varepsilon) \frac{\partial}{\partial y_{j}} \left\{ \Psi_{\varepsilon, k}(y) - P_{k}(y) - \varepsilon \chi_{k}(y/\varepsilon) \right\} \cdot \frac{\partial^{2} u_{0}}{\partial y_{i} \partial y_{k}} dy$$

$$(4.19)$$

and

$$w_{\varepsilon}^{(2)}(x) = \varepsilon \int_{\partial \widetilde{D}} \widetilde{N}_{\varepsilon}(x, y) \cdot n_{i}(y) F_{jik}(y/\varepsilon) \cdot \frac{\partial^{2} u_{0}}{\partial y_{j} \partial y_{k}} d\sigma(y)$$

$$+ \int_{\partial \widetilde{D}} \widetilde{N}_{\varepsilon}(x, y) \cdot \left\{ \frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} - \frac{\partial u_{0}}{\partial \nu_{0}} \right\} d\sigma(y)$$

$$(4.20)$$

(we have supressed all superscripts for notational simplicity).

To estimate $w_{\varepsilon}^{(2)}$ in D_1 , we note that $\mathcal{L}_{\varepsilon}(w_{\varepsilon}^{(2)}) = 0$ in \widetilde{D} and

$$\frac{\partial}{\partial \nu_{\varepsilon}} \left\{ w_{\varepsilon}^{(2)} \right\} = \varepsilon n_i F_{jik}(x/\varepsilon) \frac{\partial^2 u_0}{\partial x_j \partial x_k} - \frac{\varepsilon}{|\partial \widetilde{D}|} \int_{\partial \widetilde{D}} n_i F_{jik}(x/\varepsilon) \frac{\partial^2 u_0}{\partial x_j \partial x_k} d\sigma \qquad \text{on } \Delta_2,$$

where we have used $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\partial u_0}{\partial \nu_0}$ on Δ_2 . Since $\|\frac{\partial}{\partial \nu_{\varepsilon}} \{w_{\varepsilon}^{(2)}\}\|_{L^{\infty}(\Delta_2)} \leq C\varepsilon \|\nabla^2 u_0\|_{L^{\infty}(D_3)}$ and

$$\left\| \frac{\partial}{\partial \nu_{\varepsilon}} \left\{ w_{\varepsilon}^{(2)} \right\} \right\|_{C^{0,\rho}(\Delta_2)} \le C \varepsilon^{1-\rho} \| \nabla^2 u_0 \|_{L^{\infty}(D_2)} + C \varepsilon \| \nabla^2 u_0 \|_{C^{0,\rho}(D_2)},$$

it follows from the boundary Lipschitz estimate in [21, Theorem 7.1] that

$$\|\nabla w_{\varepsilon}^{(2)}\|_{L^{\infty}(D_{1})} \leq C\varepsilon^{1-\rho}\|\nabla^{2}u_{0}\|_{L^{\infty}(D_{2})} + C\varepsilon\|\nabla^{2}u_{0}\|_{C^{0,\rho}(D_{2})} + C\int_{D_{2}}|w_{\varepsilon}^{(2)} - c|$$

for any constant c. This leads to

$$\|\nabla w_{\varepsilon}\|_{L^{\infty}(D_{1})} \leq \|\nabla w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{1})} + \|\nabla w_{\varepsilon}^{(2)}\|_{L^{\infty}(D_{1})}$$

$$\leq C \int_{D_{2}} |w_{\varepsilon}| \, dx + C \|\nabla w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{2})} + C\varepsilon^{1-\rho} \|\nabla^{2}u_{0}\|_{L^{\infty}(D_{3})}$$

$$+ C\varepsilon \|\nabla^{2}u_{0}\|_{C^{0,\rho}(D_{3})}.$$

Since $|\Psi_{\varepsilon,j}^{\beta} - P_j^{\beta}| \le C\varepsilon \ln[\varepsilon^{-1}M + 2]$ by Proposition 2.3, we obtain

$$\|\frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \{\Psi_{\varepsilon,j}^{\alpha\beta}\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}}\|_{L^{\infty}(D_{1})}$$

$$\leq C \int_{D_{2}} |u_{\varepsilon} - u_{0}| + C\varepsilon \ln[\varepsilon^{-1}M + 2] \|\nabla u_{0}\|_{L^{\infty}(D_{3})} + C\varepsilon \ln[\varepsilon^{-1}M + 2] \|\nabla^{2}u_{0}\|_{L^{\infty}(D_{3})}$$

$$+ C\varepsilon^{1-\rho} \|\nabla^{2}u_{0}\|_{L^{\infty}(D_{3})} + C\varepsilon \|\nabla^{2}u_{0}\|_{C^{0,\rho}(D_{3})} + C\|\nabla w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_{2})}.$$

$$(4.21)$$

It remains to estimate $\nabla w_{\varepsilon}^{(1)}$ on D_2 . The first two integrals in the right hand side of (4.19) may be handled by applying Lemma 4.6 on \widetilde{D} . Indeed, let

$$f(x) = -\varepsilon F_{jik}(x/\varepsilon) \cdot \frac{\partial^2 u_0}{\partial x_i \partial x_k} - a_{ij}(x/\varepsilon) \{ \Psi_{\varepsilon,k}(x) - P_k(x) \} \frac{\partial^2 u_0}{\partial x_i \partial x_k}.$$

Note that $||f||_{L^{\infty}(\widetilde{D})} \leq C\varepsilon \ln[\varepsilon^{-1}M + 2]||\nabla^2 u_0||_{L^{\infty}(D_3)}$ and

$$H_{\varepsilon,\rho}(f) \le C\varepsilon^{1-\rho} \ln[\varepsilon^{-1}M+2] \|\nabla^2 u\|_{L^{\infty}(D_3)} + C\varepsilon \ln[\varepsilon^{-1}M+2] \|\nabla^2 u_0\|_{C^{0,\rho}(D_3)}.$$

It follows by Lemma 4.6 that the first two integrals in the right hand side of (4.19) are bounded by

$$C\varepsilon \ln[\varepsilon^{-1}M + 2] \{ \varepsilon^{-\rho} \| \nabla^2 u_0 \|_{L^{\infty}(D_3)} + \| \nabla^2 u_0 \|_{C^{0,\rho}(D_3)} \}.$$

Finally, the third integral in (4.19) is bounded by

$$C\|\nabla^2 u_0\|_{L^{\infty}(D_3)} \int_{\widetilde{D}} \frac{|\nabla_y \{\Psi_{\varepsilon,k}(y) - P_y(y) - \varepsilon \chi_k(y/\varepsilon)\}|}{|x - y|^{d-1}} dy.$$

$$(4.22)$$

Using that $|\nabla_y \{\Psi_{\varepsilon,k}(y) - P_k(y) - \varepsilon \chi_k(y/\varepsilon)\}| \le C \min(1, \varepsilon[\operatorname{dist}(y, \partial\Omega)]^{-1})$, one may show that the integral in (4.22) is bounded by $C\varepsilon[\ln(\varepsilon^{-1}+2)]^2$. Thus, we have proved that

$$\|w_{\varepsilon}^{(1)}\|_{L^{\infty}(D_2)} \le C\varepsilon^{1-\rho} \ln[\varepsilon^{-1}M+2] \|\nabla^2 u_0\|_{L^{\infty}(D_3)} + C\varepsilon \ln[\varepsilon^{-1}M+2] \|\nabla^2 u_0\|_{C^{0,\rho}(D_3)}.$$

This, together with (4.21), yields the desired estimate.

We are now in a position to give the proof of estimate (1.10).

Theorem 4.8. Suppose that A(y) satisfies conditions (1.2), (1.3) and (1.4). Also assume that $A^* = A$. Let Ω be a bounded $C^{2,\eta}$ domain for some $\eta \in (0,1)$. Then

$$\left| \frac{\partial}{\partial x_i} \left\{ N_{\varepsilon}^{\alpha\beta}(x,y) \right\} - \frac{\partial}{\partial x_i} \left\{ \Psi_{\varepsilon,j}^{\alpha\gamma}(x) \right\} \cdot \frac{\partial}{\partial x_i} \left\{ N_0^{\gamma\beta}(x,y) \right\} \right| \le \frac{C \varepsilon^{1-\rho} \ln[\varepsilon^{-1}M + 2]}{|x - y|^{d-\rho}} \tag{4.23}$$

for any $x, y \in \Omega$ and $\rho \in (0, 1)$, where C depends only on d, m, μ , λ , τ , ρ and Ω .

Proof. Since $|\nabla_x N_{\varepsilon}(x,y)| \leq C|x-y|^{1-d}$ and $|\nabla \Psi_{\varepsilon,j}^{\beta}| \leq C$, we may assume that $\varepsilon < |x-y|$ and ρ is small. Fix $x_0, y_0 \in \Omega$, $1 \leq \gamma \leq d$ and let $r = |x_0 - y_0|/8$. Let $u_{\varepsilon}^{\alpha}(x) = N_{\varepsilon}^{\alpha\gamma}(x,y_0)$ and $u_0^{\alpha}(x) = N_0^{\alpha\gamma}(x,y_0)$. Observe that

$$\begin{cases} \mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0}) = 0 & \text{in } D(x_{0}, r), \\ \left(\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}}\right)^{\alpha} = \left(\frac{\partial u_{0}}{\partial \nu_{0}}\right)^{\alpha} = -|\partial \Omega|^{-1} \delta^{\alpha \gamma} & \text{on } \Delta(x_{0}, r). \end{cases}$$

Also, note that $\|\nabla u_0\|_{L^{\infty}(D(x_0,r))} \leq Cr^{1-d}$, $\|\nabla^2 u_0\|_{L^{\infty}(D(x_0,r))} \leq Cr^{-d}$ and $\|\nabla^2 u_0\|_{C^{0,\rho}(D(x_0,r))} \leq Cr^{-d-\rho}$. Furthermore, it follows from Theorem 4.3 that

$$||u_{\varepsilon} - u_0||_{L^{\infty}(D(x_0,r))} \le C\varepsilon r^{1-d} \ln[\varepsilon^{-1}r + 2].$$

Thus, by Lemma 4.7, we obtain

$$\left\| \frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Psi_{\varepsilon,j}^{\alpha\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{i}} \right\|_{L^{\infty}(D(x_{0},r/3))} \leq C \varepsilon^{1-\rho} r^{\rho-d} \ln[\varepsilon^{-1}M + 2].$$

This completes the proof.

Remark 4.9. Under the symmetry condition $A^* = A$, one has $N_{\varepsilon}^{\alpha\beta}(x,y) = N_{\varepsilon}^{\beta\alpha}(y,x)$. Thus, it follows from (4.23) that for any $x, y \in \Omega$,

$$\left| \frac{\partial}{\partial y_i} \left\{ N_{\varepsilon}^{\alpha\beta}(x,y) \right\} - \frac{\partial}{\partial y_i} \left\{ \Psi_{\varepsilon,\ell}^{\beta\sigma}(y) \right\} \cdot \frac{\partial}{\partial y_\ell} \left\{ N_0^{\alpha\sigma}(x,y) \right\} \right| \le \frac{C \varepsilon^{1-\rho} \ln[\varepsilon^{-1}M + 2]}{|x - y|^{d-\rho}}. \tag{4.24}$$

Fix β and j. Let

$$u_{\varepsilon}^{\alpha}(x) = \frac{\partial}{\partial y_{i}} \left\{ N_{\varepsilon}^{\alpha\beta}(x,y) \right\} \text{ and } u_{0}^{\alpha}(x) = \frac{\partial}{\partial y_{i}} \left\{ \Psi_{\varepsilon,k}^{\beta\sigma}(y) \right\} \cdot \frac{\partial}{\partial y_{k}} \left\{ N_{0}^{\alpha\sigma}(x,y) \right\}.$$

Note that $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = \mathcal{L}_{0}(u_{0}) = 0$ in $\Omega \setminus \{y\}$ and $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = \frac{\partial u_{0}}{\partial \nu_{0}} = 0$ on $\partial \Omega$. We may use Lemma 4.7 and estimate (4.24) to deduce that if Ω is $C^{3,\eta}$ for some $\eta \in (0,1)$,

$$\left| \frac{\partial^{2}}{\partial x_{i} \partial y_{j}} \left\{ N_{\varepsilon}^{\alpha \beta}(x, y) \right\} - \frac{\partial}{\partial x_{i}} \left\{ \Psi_{\varepsilon, k}^{\alpha \gamma}(x) \right\} \cdot \frac{\partial^{2}}{\partial x_{k} \partial y_{\ell}} \left\{ N_{0}^{\gamma \sigma}(x, y) \right\} \cdot \frac{\partial}{\partial y_{j}} \left\{ \Psi_{\varepsilon, \ell}^{\beta \sigma}(y) \right\} \right| \\
\leq \frac{C_{\rho} \varepsilon^{1-\rho} \ln[\varepsilon^{-1} M + 2]}{|x - y|^{d+1-\rho}} \tag{4.25}$$

for any $x, y \in \Omega$ and $\rho \in (0, 1)$, where C_{ρ} depends only on $d, m, \mu, \lambda, \tau, \rho$ and Ω .

As a corollary of Theorem 4.8, we obtain an $O(\varepsilon^t)$ estimate for any $t \in (0,1)$ in $W^{1,p}(\Omega)$ for solutions with Neumann boundary conditions.

Theorem 4.10. Suppose that A(y) and Ω satisfy the same conditions as in Theorem 4.8. Let $1 . For <math>\varepsilon \ge 0$ and $F \in L^p(\Omega)$ with $\int_{\Omega} F = 0$, let $u_{\varepsilon} \in W^{1,p}(\Omega)$ be the solution of the Neumann problem: $\mathcal{L}_{\varepsilon}(u_{\varepsilon}) = F$ in Ω , $\frac{\partial u_{\varepsilon}}{\partial \nu_{\varepsilon}} = 0$ and $\int_{\partial \Omega} u_{\varepsilon} = 0$. Then

$$\|u_{\varepsilon} - u_0 - \{\Psi_{\varepsilon,j}^{\beta} - P_j^{\beta}\} \frac{\partial u_0^{\beta}}{\partial x_j} \|_{W^{1,p}(\Omega)} \le C_t \varepsilon^t \|F\|_{L^p(\Omega)}$$

$$(4.26)$$

for any $t \in (0,1)$, where C_t depends only on d, m, μ , λ , τ , t, p and Ω .

Proof. Since $\|\Psi_{\varepsilon,j}^{\beta} - P_k^{\beta}\|_{L^{\infty}(\Omega)} \le C\varepsilon \ln[\varepsilon^{-1}M + 2]$ and $\|\nabla u_0\|_{L^p(\Omega)} + \|\nabla^2 u_0\|_{L^p(\Omega)} \le C\|F\|_{L^p(\Omega)}$, in view of Theorem 4.5, it suffices to prove that

$$\left\| \frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Psi_{\varepsilon, j}^{\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{i}} \right\|_{L^{p}(\Omega)} \le C_{t} \varepsilon^{t} \|F\|_{L^{p}(\Omega)}. \tag{4.27}$$

We will prove that estimate (4.27) holds for any $1 \leq p \leq \infty$. To this end, note that $u_{\varepsilon}(x) = \int_{\Omega} N_{\varepsilon}(x, y) F(y) dy$, by Theorem 4.8,

$$\left| \frac{\partial u_{\varepsilon}^{\alpha}}{\partial x_{i}} - \frac{\partial}{\partial x_{i}} \left\{ \Psi_{\varepsilon,j}^{\beta} \right\} \cdot \frac{\partial u_{0}^{\beta}}{\partial x_{j}} \right| \leq C_{\rho} \, \varepsilon^{1-\rho} \ln[\varepsilon^{-1}M + 2] \int_{\Omega} \frac{|F(y)| \, dy}{|x - y|^{d-\rho}} \tag{4.28}$$

for any $\rho \in (0,1)$. Note that if $\rho < 1-t$, then $\varepsilon^{1-\rho} \ln[\varepsilon^{-1}M+2] \leq C\varepsilon^t$. Estimate (4.27) follows easily from (4.28).

5 Leibniz Rules for the Dirichlet-to-Neumann map

Let $\Lambda: H^{1/2}(\partial\Omega) \to H^{-1/2}(\partial\Omega)$ denote the Dirichlet-to-Neumann map associated with $\mathcal{L} = -\Delta$. It is known that the estimate $\|\Lambda(f)\|_{L^p(\partial\Omega)} \leq C_p \|f\|_{W^{1,p}(\partial\Omega)}$ holds for $1 if <math>\Omega$ is C^1 [12], and for $1 if <math>\Omega$ is Lipschitz [19, 27, 9]. The goal of this section is to prove the Leibniz estimates that were used in subsection 3.3. We remark that our basic line of argument extends to the general case $\mathcal{L} = -\text{div}(A(x)\nabla)$. Indeed, Theorem 5.2 continues to hold for if A(x) satisfies conditions (1.2), $\|\nabla A\|_{\infty} \leq C$ and $A^* = A$. Under the same conditions on A, Theorem 5.1 also holds in the scalar case (m = 1). If m > 1, estimate (5.1) holds for C^1 domains. The details will be given elsewhere.

Theorem 5.1. Let $\mathcal{L} = -\Delta$ and Ω be a bounded Lipschitz domain. Then

$$\|\Lambda(fg) - f\Lambda(g)\|_{L^2(\partial\Omega)} \le C\|f\|_{H^1(\partial\Omega)}\|g\|_{L^\infty(\partial\Omega)} \tag{5.1}$$

for any $f \in H^1(\partial\Omega)$ and $g \in L^{\infty}(\partial\Omega)$.

Proof. Let w be the solution to the L^2 Dirichlet problem with data fg; i.e., $\mathcal{L}(w) = 0$ in Ω , w = fg on $\partial\Omega$ and $(w)^* \in L^2(\partial\Omega)$, where $(w)^*$ denotes the nontangential maximal function of w. Let u and v denote the solutions of the L^2 Dirichlet problem with boundary data f and g respectively. Then

$$\Lambda(fg) - f\Lambda(g) - g\Lambda(f) = \frac{\partial}{\partial \nu} \{w - uv\}.$$

To estimate $\partial(w-uv)/\partial\nu$ in $L^2(\partial\Omega)$, let h be a function such that $\mathcal{L}(h)=0$ in Ω and $(h)^* \in L^2(\partial\Omega)$. Since w-uv=0 on $\partial\Omega$ and

$$\mathcal{L}(w - uv) = 2 \frac{\partial u}{\partial x_j} \frac{\partial v}{\partial x_j} = 2 \frac{\partial}{\partial x_j} \left\{ \frac{\partial u}{\partial x_j} \cdot v \right\},\,$$

by the Green's formula, we obtain

$$\int_{\partial\Omega} \frac{\partial}{\partial\nu} \{w - uv\} \cdot h = 2 \int_{\Omega} \frac{\partial u}{\partial x_j} \cdot v \cdot \frac{\partial h}{\partial x_j} - 2 \int_{\partial\Omega} \Lambda(f) \cdot g \cdot h. \tag{5.2}$$

Thus, by duality, it suffices to prove the following estimate

$$\left| \int_{\Omega} \frac{\partial u}{\partial x_j} \cdot v \cdot \frac{\partial h}{\partial x_j} \right| \le C \|f\|_{H^1(\partial\Omega)} \|g\|_{L^{\infty}(\partial\Omega)} \|h\|_{L^2(\partial\Omega)}. \tag{5.3}$$

The trilinear estimate (5.3) on Lipschitz domains is a consequence of a bilinear estimate due to B. Dahlberg [8] (also see [18] for related work). Indeed, since $\Delta(h) = 0$ in Ω and $(h)^* \in L^2(\partial\Omega)$, it follows from estimate (1.7) in [8] that the left hand side of (5.3) is bounded by

$$C||h||_{L^{2}(\partial\Omega)} \left\{ \int_{\partial\Omega} |((\nabla u)v)^{*}|^{2} d\sigma + \int_{\Omega} |\nabla^{2}u|^{2}|v|^{2} \delta(x) dx + \int_{\Omega} |\nabla u|^{2}|\nabla v|^{2} \delta(x) dx \right\}^{1/2}, \quad (5.4)$$

where $\delta(x) = \operatorname{dist}(x, \partial\Omega)$. The first two integals in (5.4) are bounded by $C \|f\|_{H^1(\partial\Omega)}^2 \|g\|_{L^{\infty}(\partial\Omega)}^2$. This follows from estimates $\|(\nabla u)^*\|_{L^2(\partial\Omega)} \leq C \|f\|_{H^1(\partial\Omega)}$ and $\int_{\Omega} |\nabla^2 u|^2 \delta(x) dx \leq C \|f\|_{H^1(\partial\Omega)}^2$ [19, 7] as well as the maximum principle $\|v\|_{L^{\infty}(\Omega)} \leq \|g\|_{L^{\infty}(\partial\Omega)}$.

Finally, to handle the last integral in (5.4), we use the fact that $|\nabla v(x)|^2 \delta(x) dx$ is a Carleson measure whose norm is less than $C \|g\|_{BMO(\partial\Omega)}^2$ [13]. This implies that the integral is bounded by $C \|(\nabla u)^*\|_{L^2(\partial\Omega)}^2 \|g\|_{L^\infty(\partial\Omega)}^2$. Estimate (5.3) now follows.

If Ω is smooth, then Λ is a pseudo-differential operator of order one. In this case the L^p boundedness of the commutator $[\Lambda, x_i]$ in the next theorem is well known.

Theorem 5.2. Let $\mathcal{L} = -\Delta$ and Ω be a bounded C^1 domain. Then

$$\|\Lambda(fx_i) - x_i\Lambda(f)\|_{L^p(\partial\Omega)} \le C_p \|f\|_{L^p(\partial\Omega)} \tag{5.5}$$

for any $1 and <math>f \in L^p(\partial\Omega)$. If Ω is Lipschitz, the estimate (5.5) holds for p = 2.

Proof. Let Ω be a bounded C^1 domain. Let w and u be harmonic functions in Ω with boundary data fx_i and f respectively such that $(w)^*, (u)^* \in L^p(\partial\Omega)$. Then $\Lambda(fx_i) - f\Lambda(x_i) - x_i\Lambda(f) = \partial\{w - x_iu\}/\partial\nu$. Let h be a harmonic function in Ω such that $(h)^* \in L^{p'}(\partial\Omega)$. Since $w - x_iu = 0$ on $\partial\Omega$ and $\mathcal{L}(w - x_iu) = 2\frac{\partial u}{\partial x_i}$,

$$\left| \int_{\partial\Omega} \frac{\partial}{\partial \nu} \left\{ w - x_i u \right\} \cdot h \right| = 2 \left| \int_{\Omega} \frac{\partial u}{\partial x_i} \cdot h \right| \le C \|f\|_{L^p(\partial\Omega)} \|h\|_{L^{p'}(\partial\Omega)}. \tag{5.6}$$

By duality this gives (5.5). We remark that the inequality in (5.6) was proved in [10] for p = 2 in Lipschitz domains. The same argument there also gives the inequality for $1 in <math>C^1$ domains. Thus we may conclude that estimate (5.5) holds for $1 if <math>\Omega$ is C^1 , and for p = 2 if Ω is Lipschitz.

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